
XI.

Atmospheric Muons and Neutrinos

Astroparticle Physics a.a. 2021/22

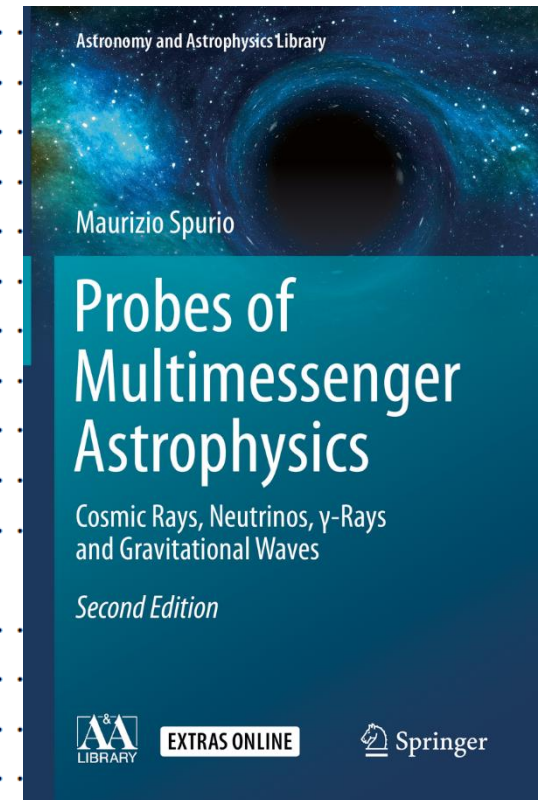
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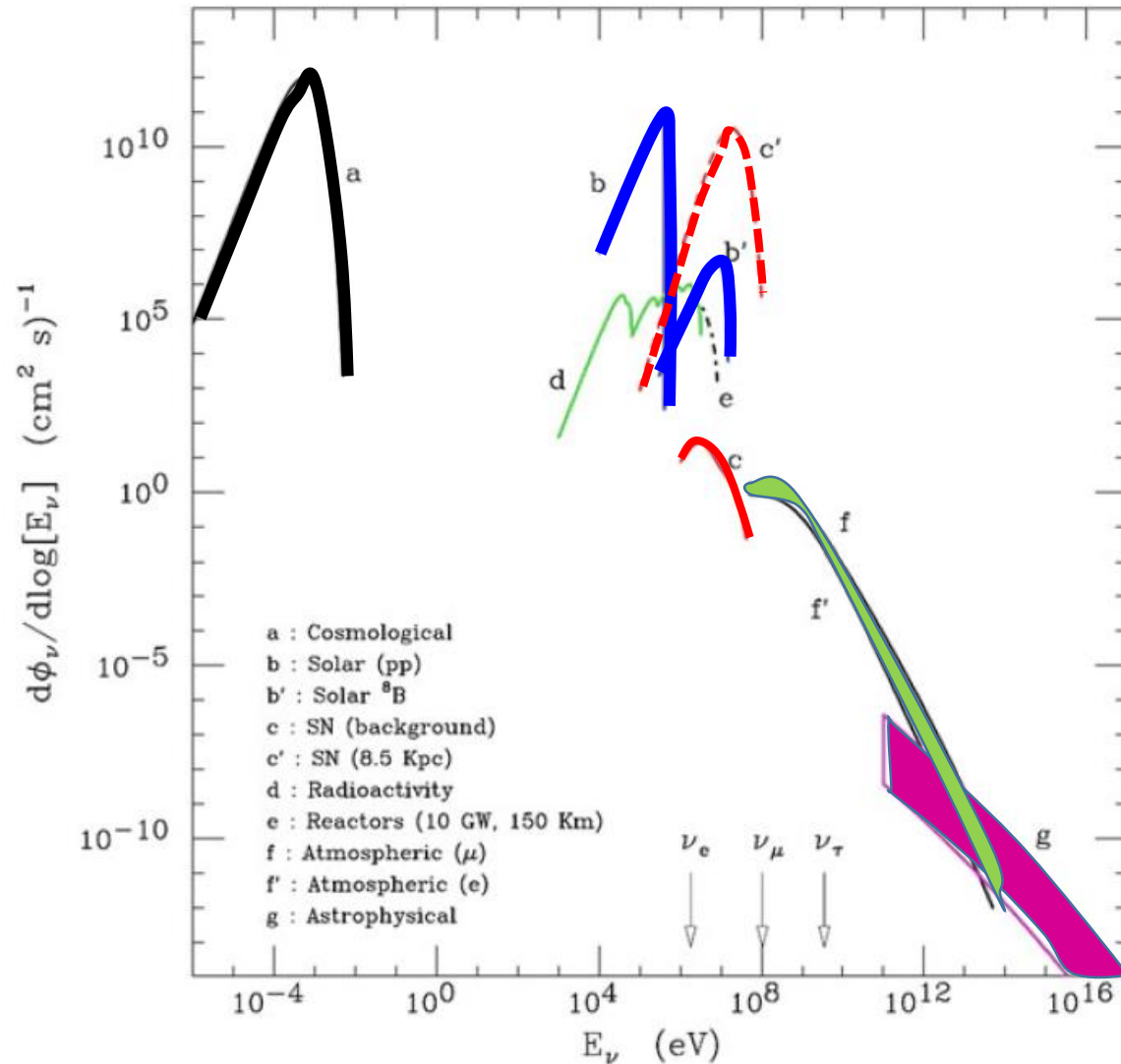
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Neutrinos from the Cosmos



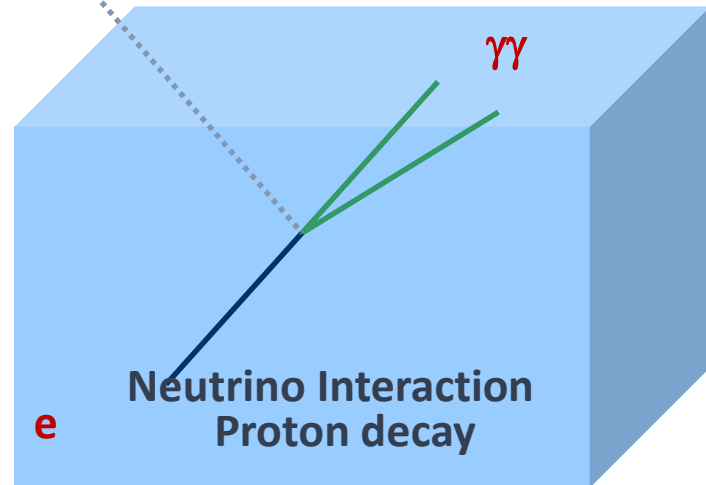
Flux of neutrinos at the surface of the Earth.

Arrow indicate the energy threshold for CC production of the charged lepton

- Big Bang neutrinos (no)
- Neutrinos from the Sun (Cap. 12)
- Neutrinos from SNe (Cap. 12)
- Atmospheric neutrinos (Here)
- High-energy cosmic neutrinos (Cap. 10)

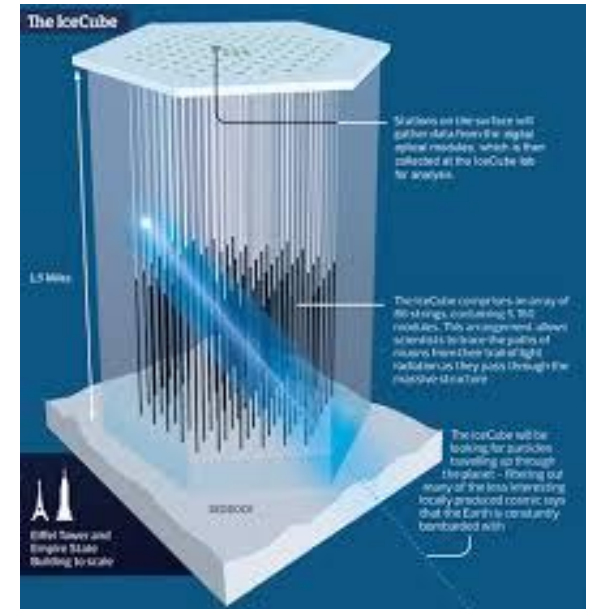
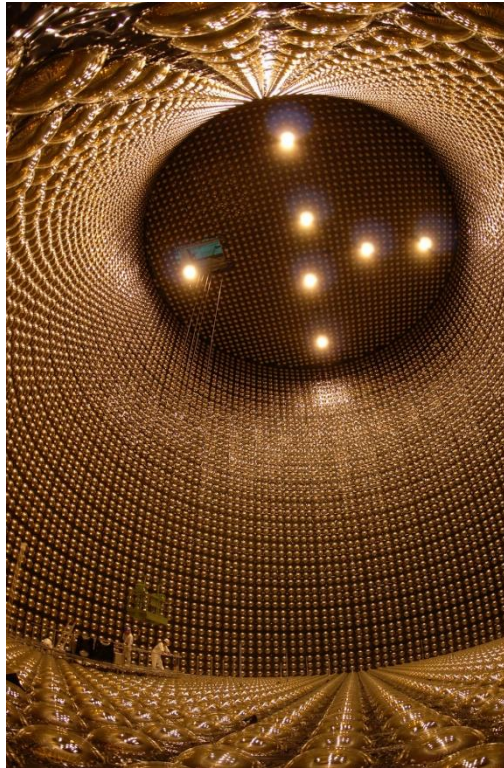
Once upon a time... GUT

- In the '80 Grand Unified Theories (**GUT**) predicted the proton decay with measurable lifetime of the order of $\tau=10^{30}$ y (20 orders of magnitude $1/H_0$)
- **QUESTION: Which is $1/H_0$?**
- The proton was thought to decay in (for instance) $p \rightarrow e^+ \pi^0$
- Detector size: 10^3 m³, and mass 1kt ($=10^{31}$ p)
- The main background for the detection of proton decay were atmospheric neutrinos interacting inside the experiment
- **QUESTION: How many p decays in 1 kt for $\tau= 10^{30}$ y, assuming 100% efficiency?**
- **Water Cerenkov** Experiments (IMB, Kamiokande)
- **Tracking calorimeters** (NUSEX, Frejus, KGF)
- **Result:** NO p decay ! But some anomalies on the neutrino measurement!

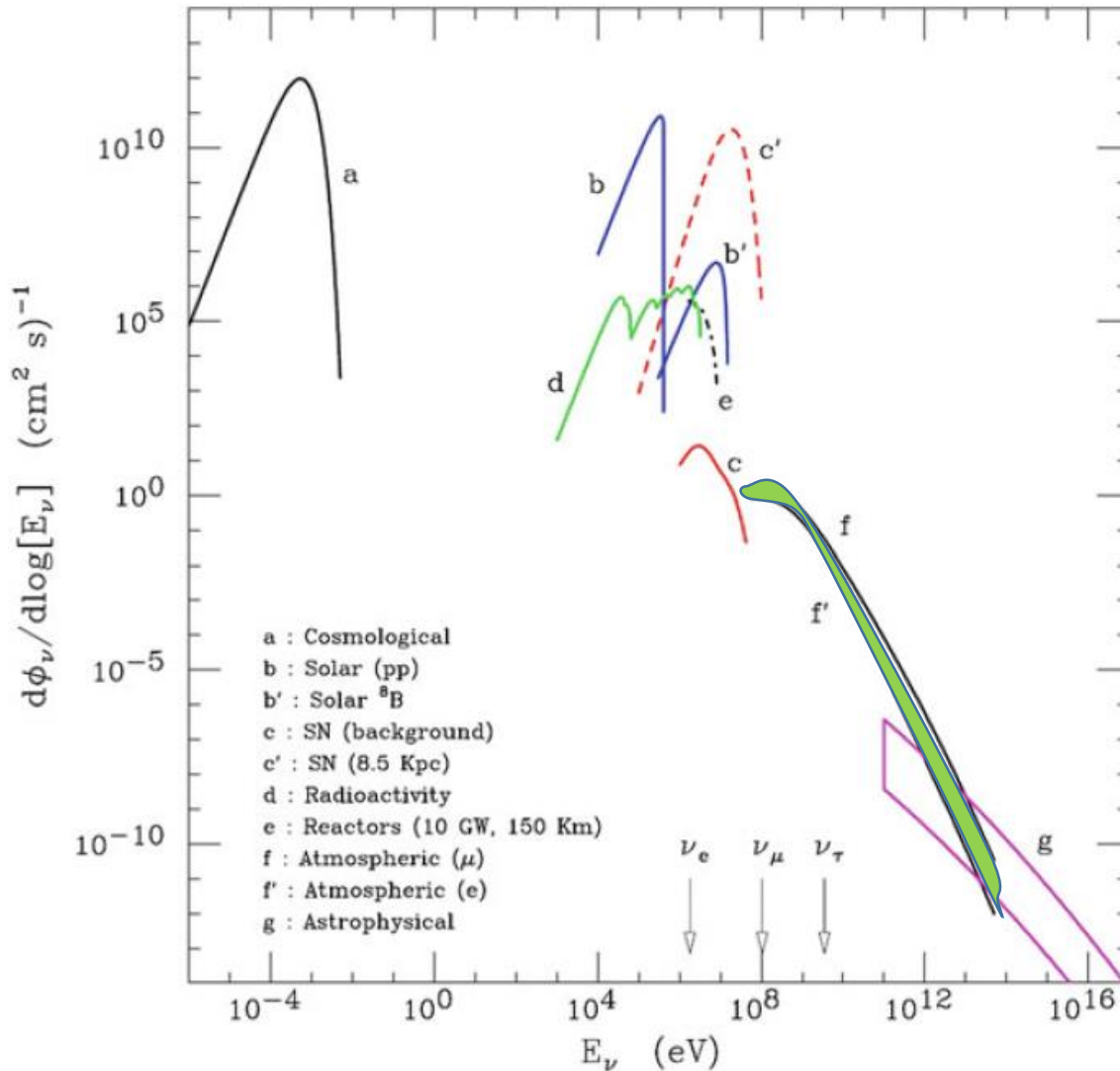


General problems for ν detectors

- Low cross section \rightarrow Large detector volume/mass
- Particle identification
- Energy/momentum measurement
- Direction measurement
- No magnetic field ($\nu = \bar{\nu}$)
- Backgrounds



Atmospheric neutrinos



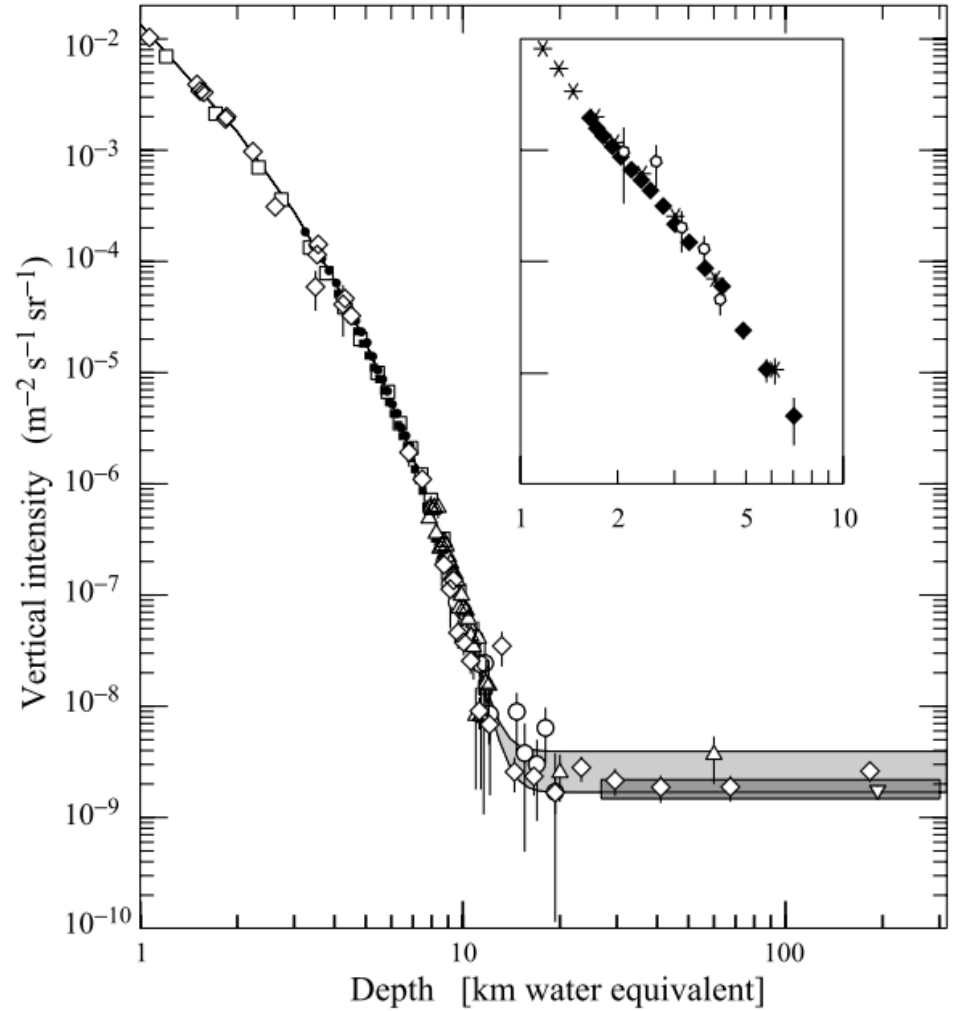
Flux of neutrinos at the surface of the Earth.

Arrow indicate the energy threshold for CC production of the charged lepton

- Big Bang neutrinos
- Neutrinos from the Sun
- Neutrinos from SNe
- **Atmospheric neutrinos**
- High-energy cosmic neutrinos

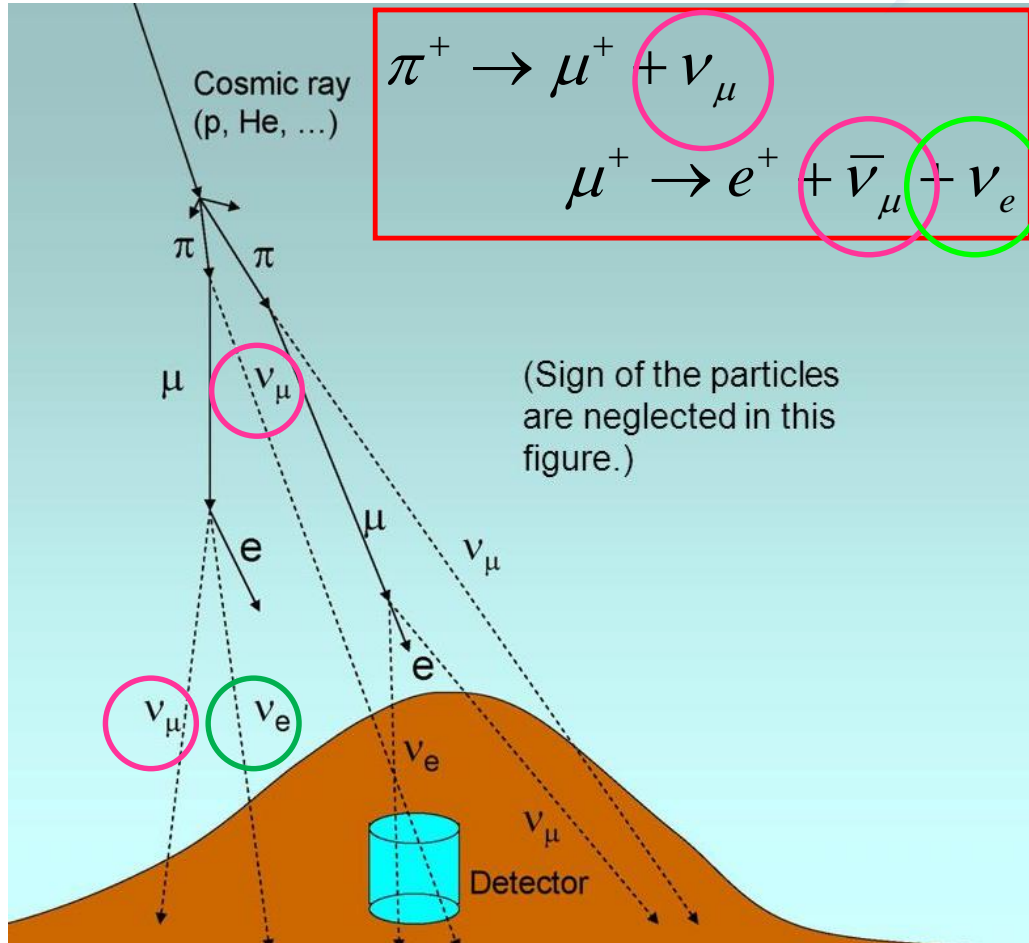
The flux of atmospheric muons vs depth

- At sea level, the muon flux is
$$\Phi \cong 100 \text{ m}^{-2} \text{ s}^{-1} \text{ sr}^{-1}$$
- Vertical muon intensity vs depth (x-axis, in km.w.e= 10^5 g cm^{-2})
- The shaded area at large depths represents neutrino-induced muons of energy $>2 \text{ GeV}$.
- [The figure inset shows the vertical intensity measured under water/ice from nu telescopes.]
- Muons cannot cross more than $\approx 15 \text{ km.w.e.}$



Recipes for the evaluation of the atmospheric flux

SKYP



Independently from the details of the computation of $\Phi_{\nu_\mu}(E)$, $\Phi_{\nu_e}(E)$, one can obtain two very robust properties:

- At energies below few GeV, the flux of ν_μ is approximately twice as large as the ν_e , i.e.:

$$\Phi(\nu_\mu) = 2\Phi(\nu_e)$$

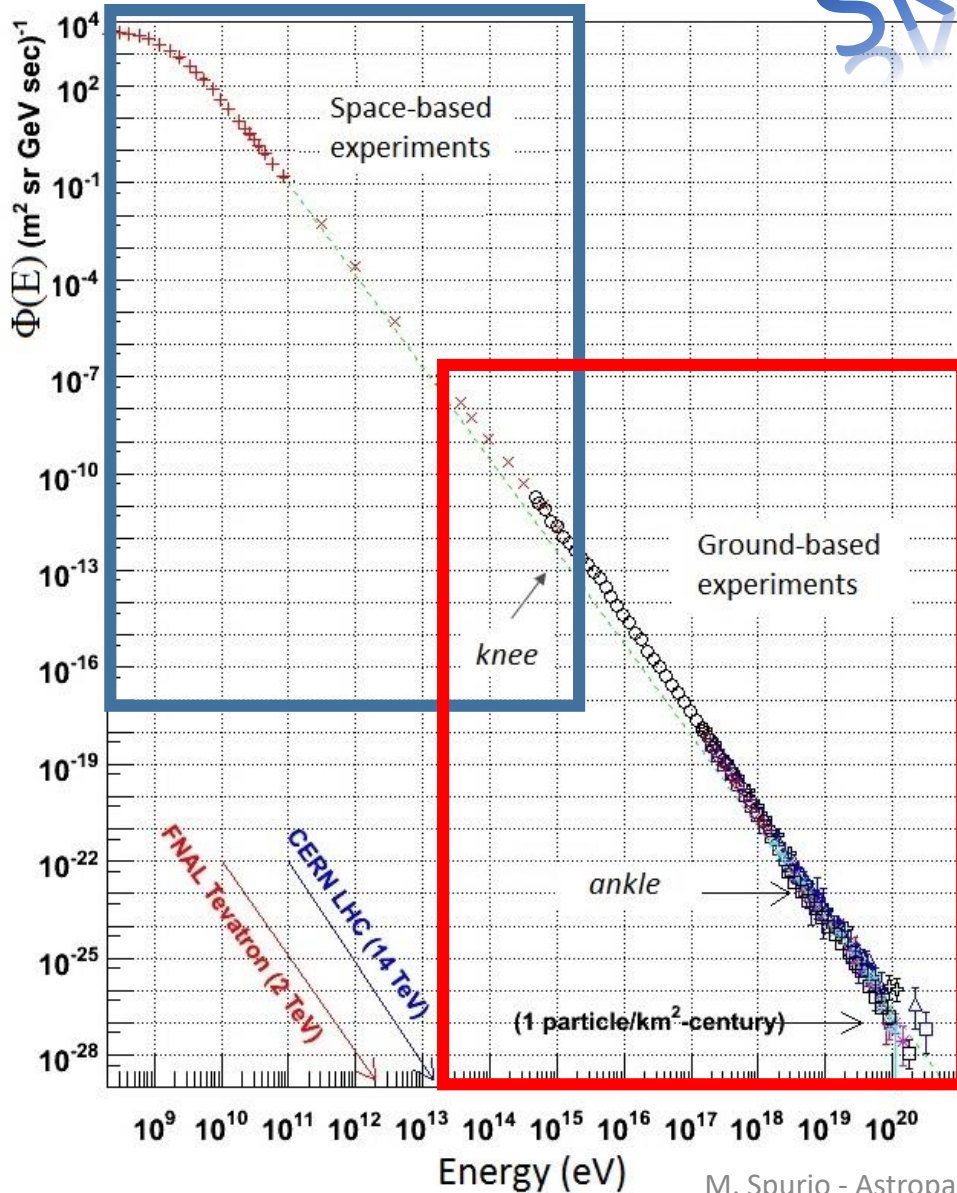
- The ν_μ , ν_e fluxes are up-down symmetric in zenith ϑ , i.e.:

$$\Phi_\nu(E_\nu, \theta) = \Phi_\nu(E_\nu, \pi - \theta)$$

Question 1: why $\Phi(\nu_\mu) = 2\Phi(\nu_e)$ not hold at higher (\gg GeV) energies?

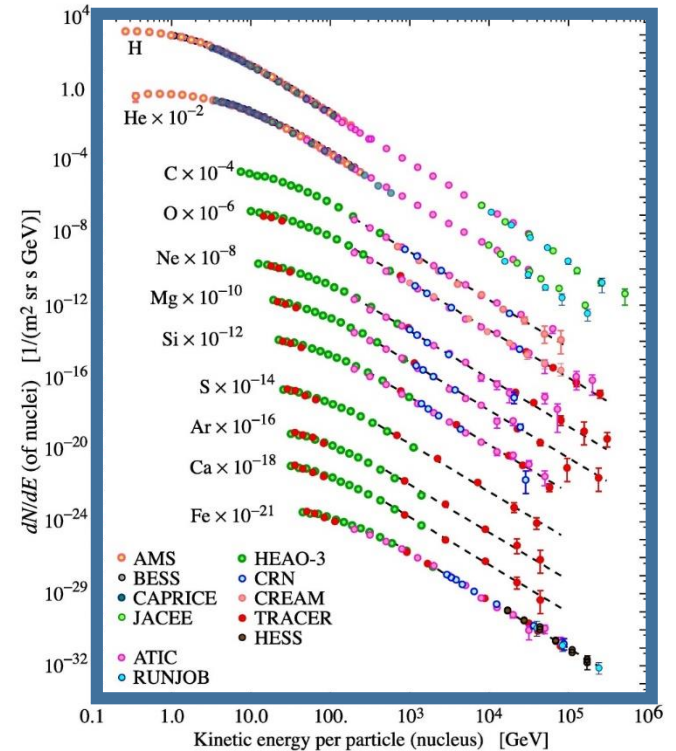
Ingredient 1) The primary CR spectrum

SKYP



Direct measurements

Indirect measurements



Ingredient 1) Cascade initiated by primary CRs

- Primary CR attenuation as function of X (g cm^{-2}) and E

$$\frac{\partial \mathcal{N}_N(E, X)}{\partial X} = -\frac{\mathcal{N}_N(E, X)}{\lambda_{IN}} + \int_E^{\infty} \frac{\mathcal{N}_N(E', X)}{\lambda_{IN}} F_{NN}(E, E') \frac{dE'}{E}$$

- Boundary condition:

$$\mathcal{N}_N(E, 0) = \Phi(E) = KE^{-\alpha}$$

- From Feynman scaling:

$$\mathcal{N}_N(E, X) = \Phi_N(E) \cdot H_N(X)$$

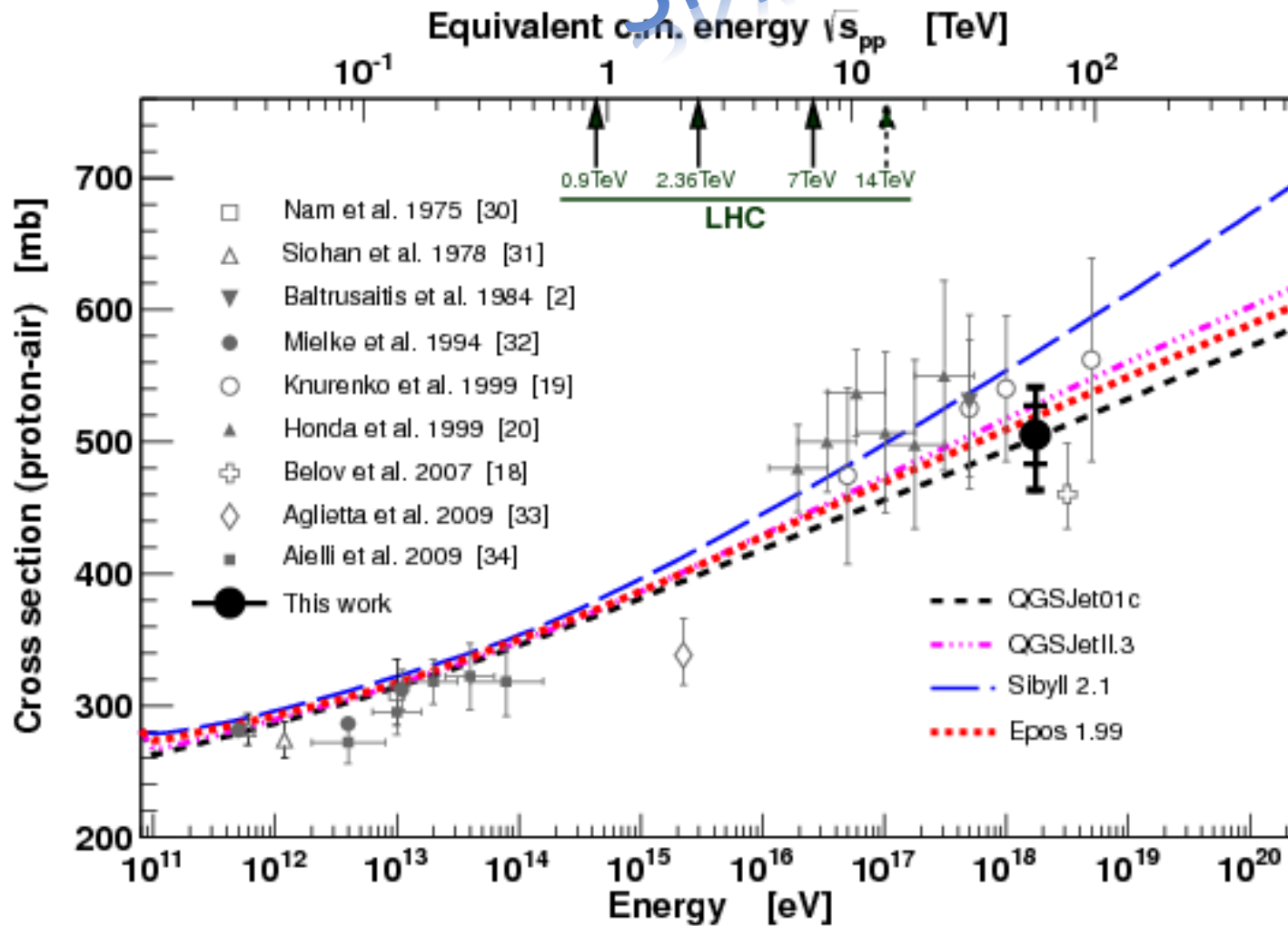
- The dependence on X depends on an effective attenuation length Λ_N :

$$H_N(X) = H_N(0) \cdot \exp\left(-\frac{X}{\Lambda_N}\right)$$

$$\Phi_N(E) = KE^{-\alpha} .$$

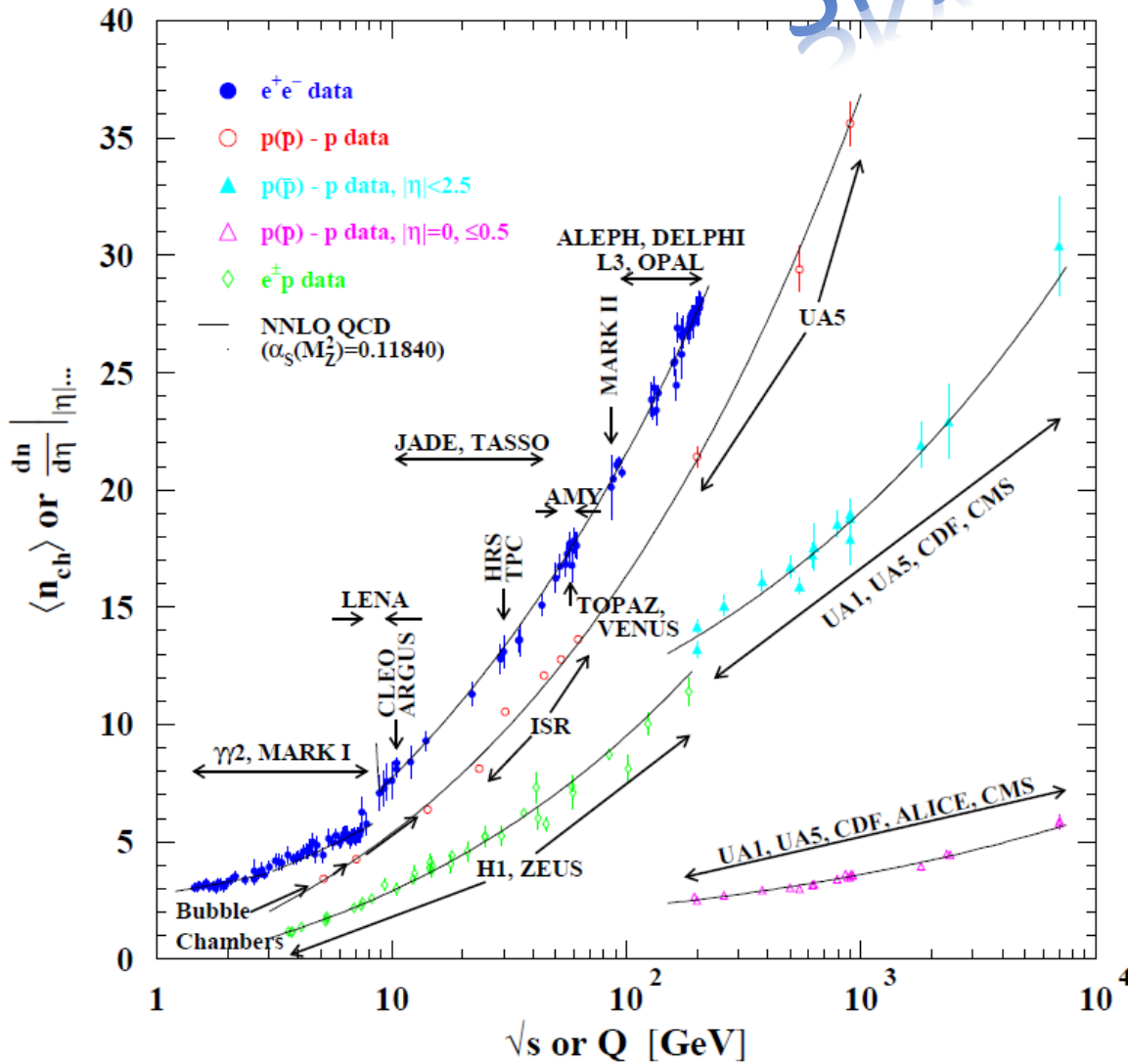
Ingredient 2) p-air cross section

SKYR



AUGER Coll. PRL 109 (2012) 062002

Ingredient 3) Secondary charged multiplicity



Average number of charged hadrons produced in pp (and $\bar{p}p$), e^+e^- , ep collisions versus center of mass energy

Ingredient 2), 3) → meson production

SKYP

- The pion propagation in atmosphere is described by:

$$\frac{\partial \mathcal{N}_\pi(E, X)}{\partial X} = - \left(\frac{1}{\lambda_{I\pi}} + \frac{1}{d_\pi} \right) \mathcal{N}_\pi(E, X) \quad (11.15)$$

$$+ \int_0^1 \frac{\mathcal{N}_\pi(\frac{E}{x^*}, X)}{\lambda_{I\pi}} \cdot F_{\pi\pi}(x^*) \cdot \frac{dx^*}{x^{*2}} + \int_0^1 \frac{\mathcal{N}_N(\frac{E}{x^*}, X)}{\lambda_{IN}} \cdot F_{N\pi}(x^*) \cdot \frac{dx^*}{x^{*2}}$$

- Competition between interaction and decay. The decay length:

$$\frac{1}{d_\pi} = \frac{1}{\Gamma c \tau_{\pi^\pm} \rho(X)} = \frac{m_\pi c^2 h_0}{E c \tau_{\pi^\pm} X \cos \theta} = \frac{\varepsilon_\pi}{E X \cos \theta}$$

- The pion decay constant $\varepsilon_\pi = 115 \text{ GeV}$.
- Pions start to increase with increasing depth X , reach a maximum and then decrease

- High-energy limit ($E \gg \varepsilon_\pi$): the decay term d_π can be neglected:

$$\mathcal{N}_\pi^{he}(E, X) = \left[\frac{Z_{N\pi}}{1 - Z_{NN}} \cdot \frac{\Lambda_\pi}{\Lambda_\pi - \Lambda_N} (e^{-X/\Lambda_\pi} - e^{-X/\Lambda_N}) \right] \cdot KE^{-\alpha}.$$

(the Z are the the *spectrum weighted moments*)

- Low-energy limit ($E \cos\vartheta \ll \varepsilon_\pi$): we neglect the term $\lambda_{I\pi}$.

$$\mathcal{N}_\pi^{le}(E, X) = \left[\frac{Z_{N\pi}}{\lambda_{IN}} \cdot e^{-X/\Lambda_N} \cdot X \right] \cdot \frac{\cos\theta}{\varepsilon_\pi} \cdot KE^{-\alpha+1}$$

- Similar equations hold for other particles with different decay constants:

Table 11.2 Lifetime τ_i and decay constants $\varepsilon_i = mc^2 h_0 / c \tau_0$ for secondary particles i produced by primary hadrons

Particles						
$i =$	μ^\pm	π^\pm	π^0	K^\pm	D^\pm	D^0
τ_i (s)	2.19×10^{-6}	2.60×10^{-8}	8.4×10^{-17}	1.24×10^{-8}	1.04×10^{-12}	4.10×10^{-13}
ε_i (GeV)	1.0	115	3.5×10^{10}	850	4.3×10^7	9.2×10^7

Ingredient 2), 3) → neutrino production

SKYP

- **Conventional** muons are produced by π and K decays:

$$\left. \begin{aligned} \pi^+(K^+) &\rightarrow \nu_\mu + \mu^+ \\ &\hookrightarrow \mu^+ \rightarrow \bar{\nu}_\mu + \nu_e + e^+ \\ \pi^-(K^-) &\rightarrow \bar{\nu}_\mu + \mu^- \\ &\hookrightarrow \mu^- \rightarrow \nu_\mu + \bar{\nu}_e + e^- \end{aligned} \right\} \begin{aligned} \frac{d\mathcal{N}_\pi}{dX} &= -\frac{\mathcal{N}_\pi}{d\pi} \\ d\mathcal{N}_\mu &= -d\mathcal{N}_\pi \end{aligned}$$

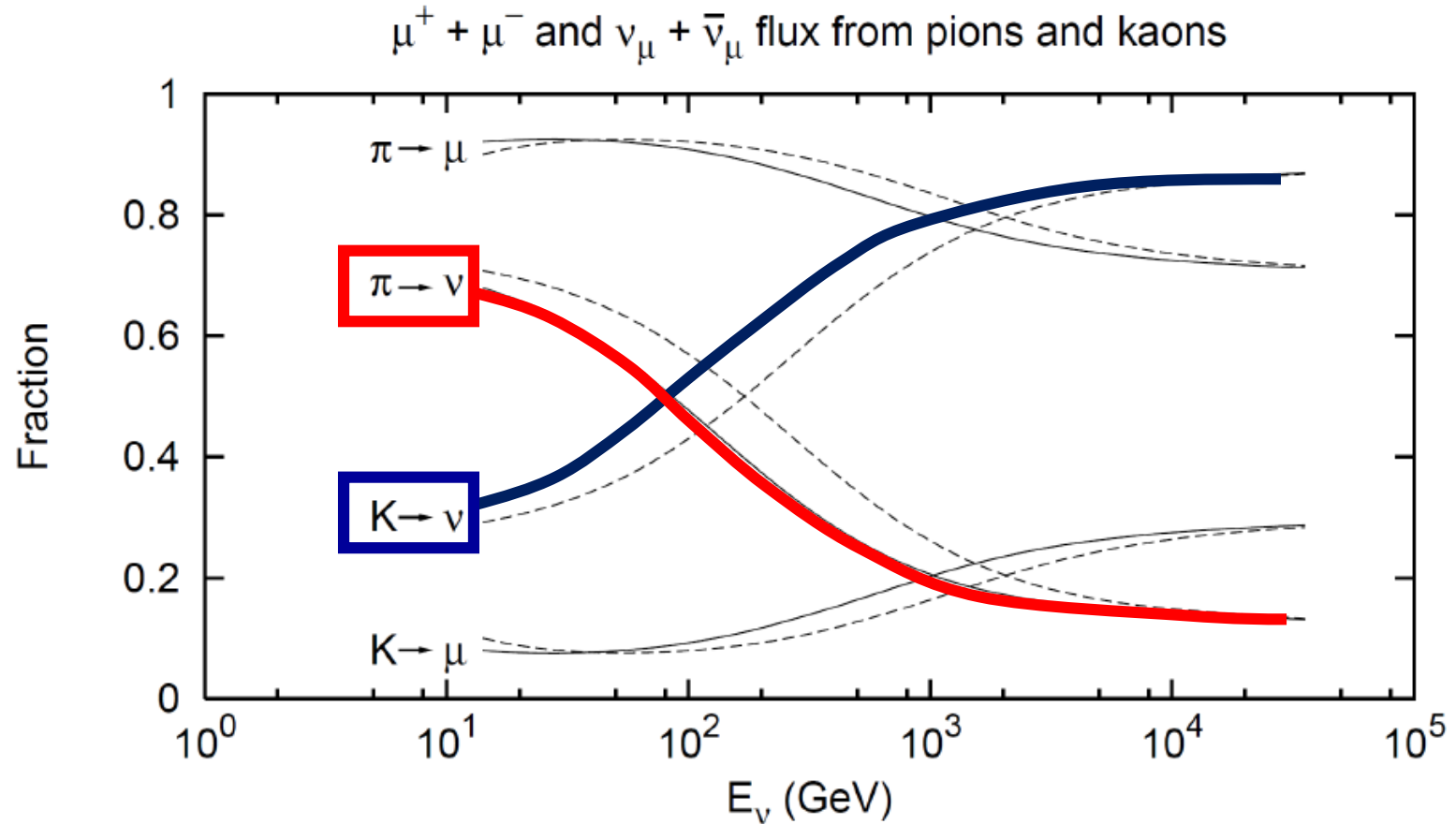
- **Prompt** muons by the decay of charmed mesons.
- The muon flux is thus described by the equation:

$$\Phi_\mu(E) = KE^{-\alpha} \left(\frac{A_\pi}{1 + \left(\frac{B_\pi E}{\varepsilon_\pi}\right) \cos \theta} + \frac{A_K}{1 + \left(\frac{B_K E}{\varepsilon_K}\right) \cos \theta} + \frac{A_c}{1 + \left(\frac{B_c E}{\varepsilon_c}\right) \cos \theta} \right)$$

- The muon neutrino flux follows similarly.

The conventional ν flux: π and K

Solid lines: vertical, dashed lines: zenith 60°

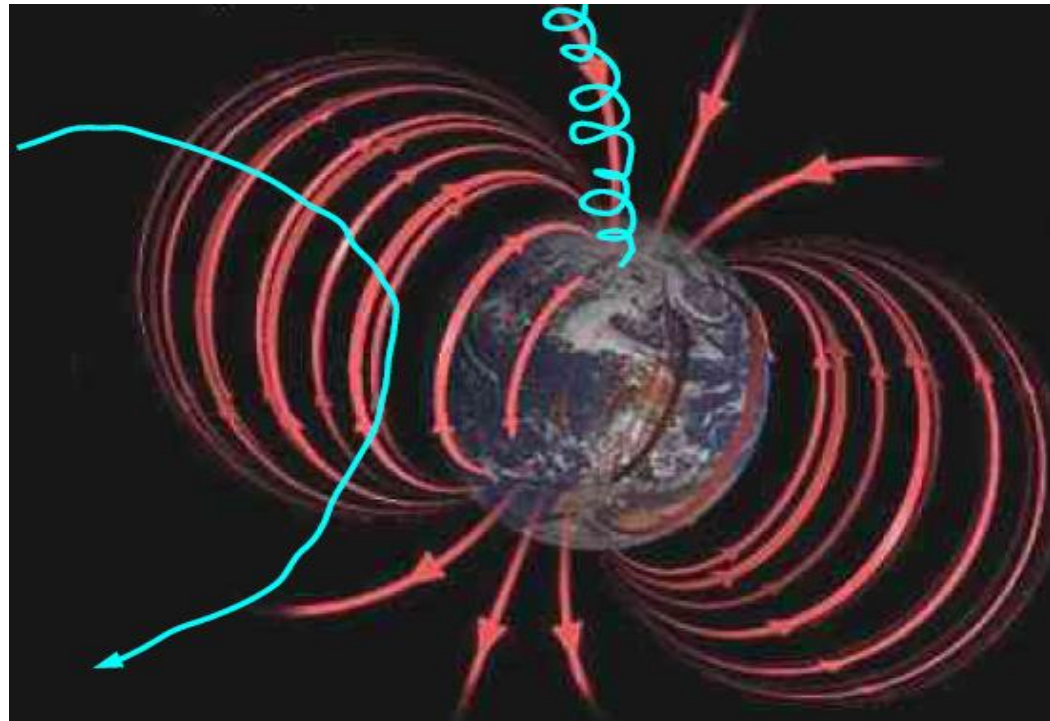


T. K. Gaisser, *Earth Planets Space* **58** (2006), 1-5.

Ingredient 4) Solar effects+geomagnetic field

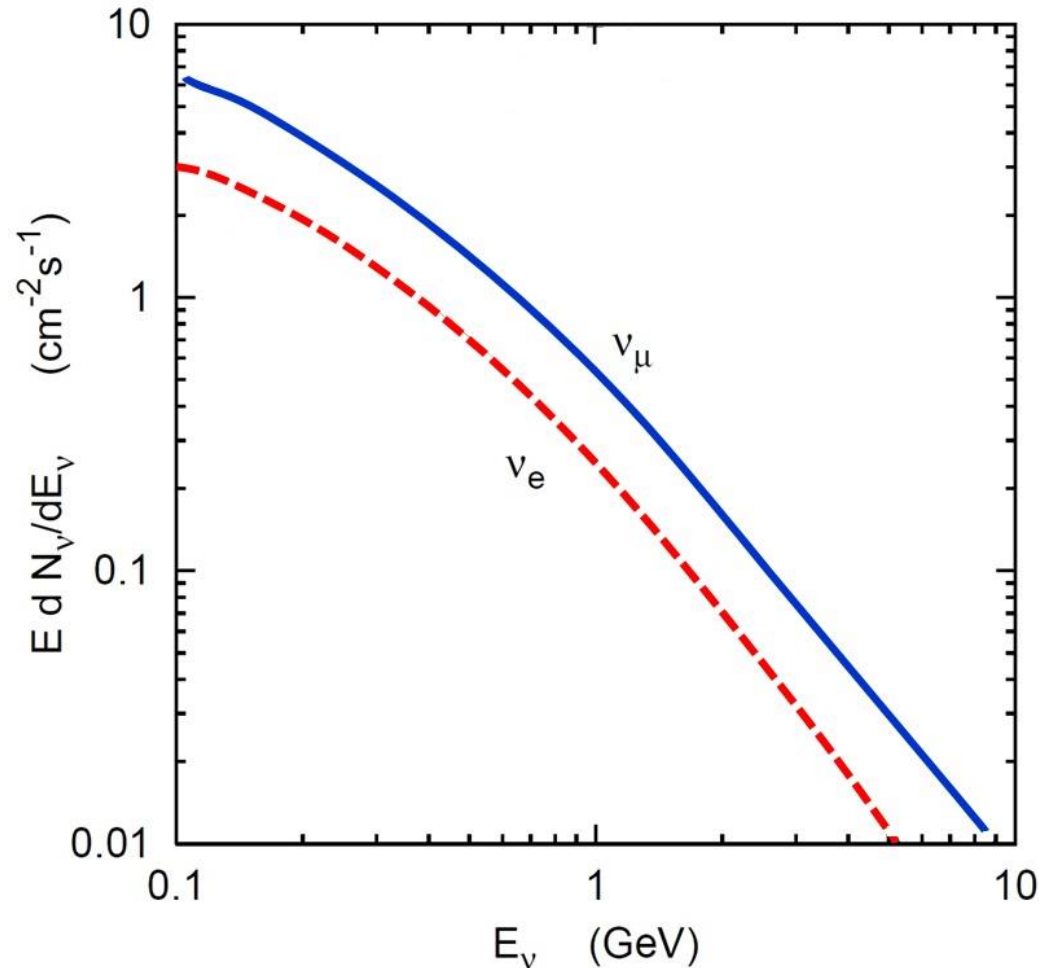
SKYP

- high precision 3D calculations,
- refined geomagnetic cut-off treatment (also geomagnetic field in atmosphere)
- elevation models of the Earth
- different atmospheric profiles
- geometry of detector effects



The conventional ν flux (Honda)

M. Honda, et al. Phys. Rev. D 92, 023004 (2015)



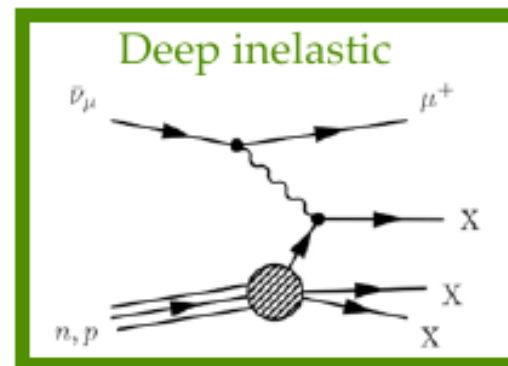
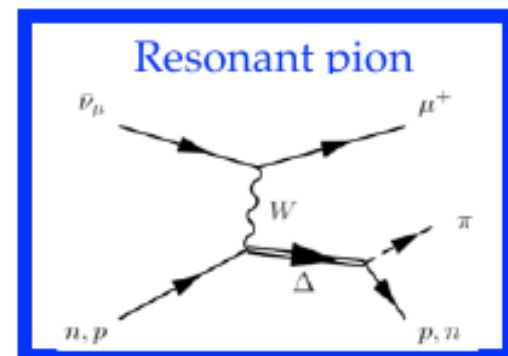
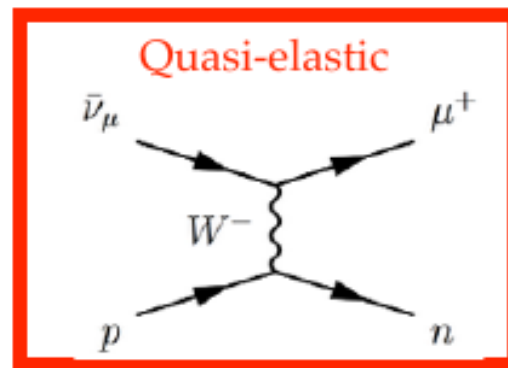
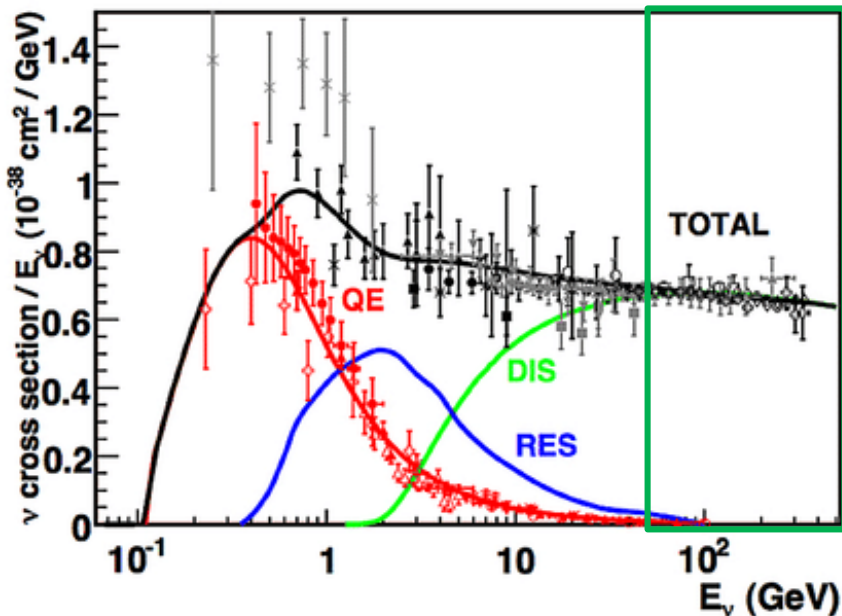
- The result of one of such calculations: the ν flux prediction (Honda) from MC simulations;
- Different models exist;

Question 2: Compute the number of ν interaction/y for $E_\nu > 1$ GeV in 1 kton detector ($\varepsilon=1$)

Question 3: Explain qualitatively the (ν_μ/ν_e) ratio

Question 4: Explain qualitatively why the $(\nu_\mu/\bar{\nu}_\mu)$ ratio increases with energy

Neutrino interactions and cross section



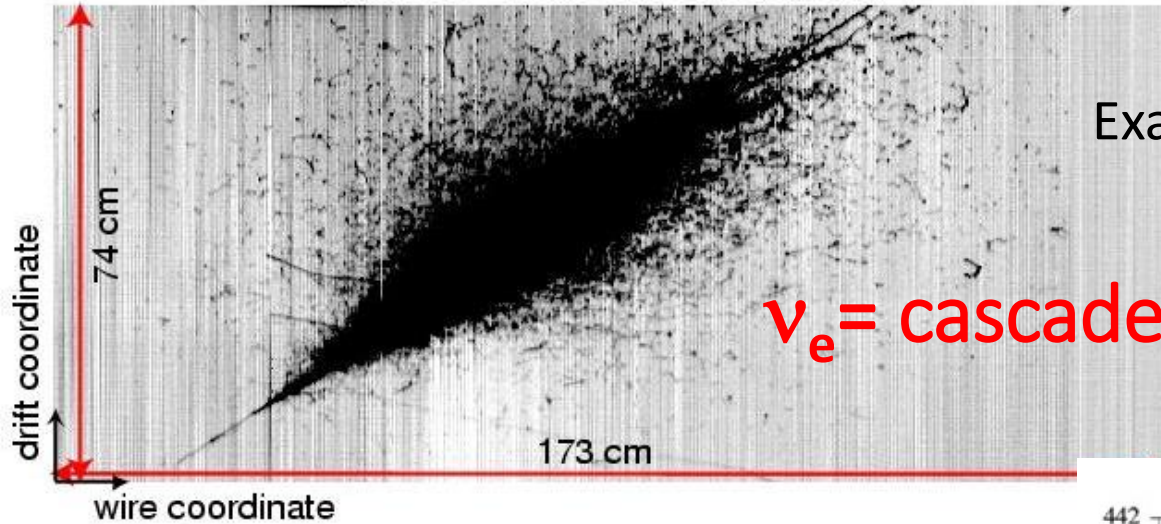
Energy threshold

- $\nu_e > 0.5 \text{ MeV}$
- $\nu_\mu > 105 \text{ MeV}$
- $\nu_\tau > 1700 \text{ MeV}$

- Different processes for the neutrino interaction on nucleons
- QE: warning to the conservation laws!
- **Question:** why we do not mention the neutrino cross-section on electrons?
- Which is the expected behavior for $\nu+e$?

Measurement of atmospheric ν 's

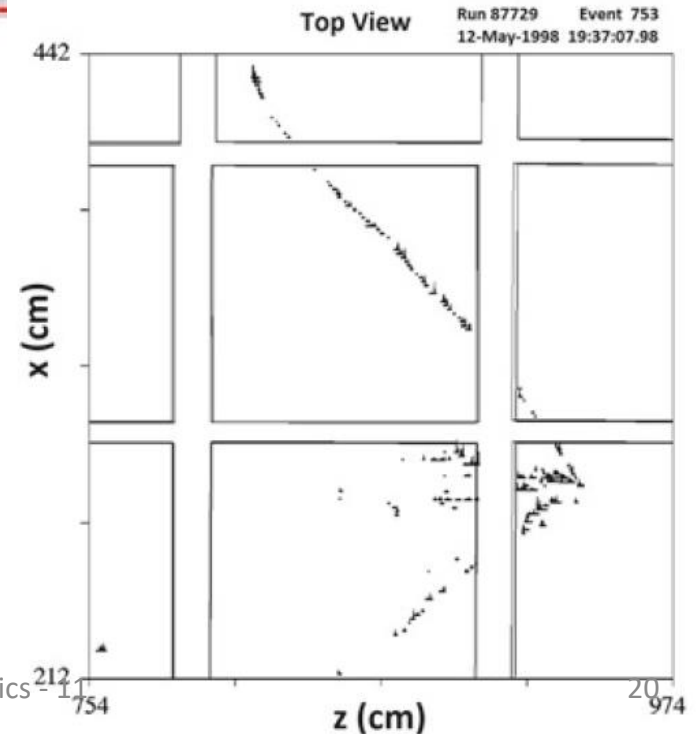
Run 308 Event 332 Collection view



Example: Icarus@ Gran Sasso

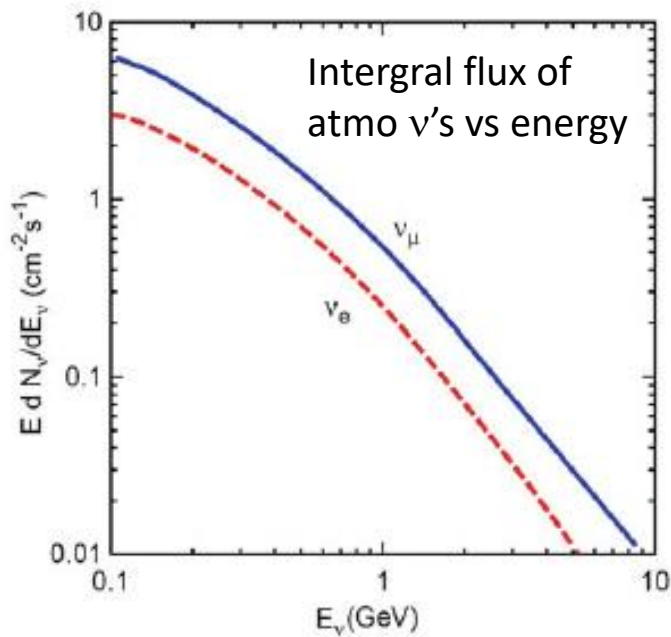
$\nu_\mu = \text{track}$

Example: Soudan II@ USA

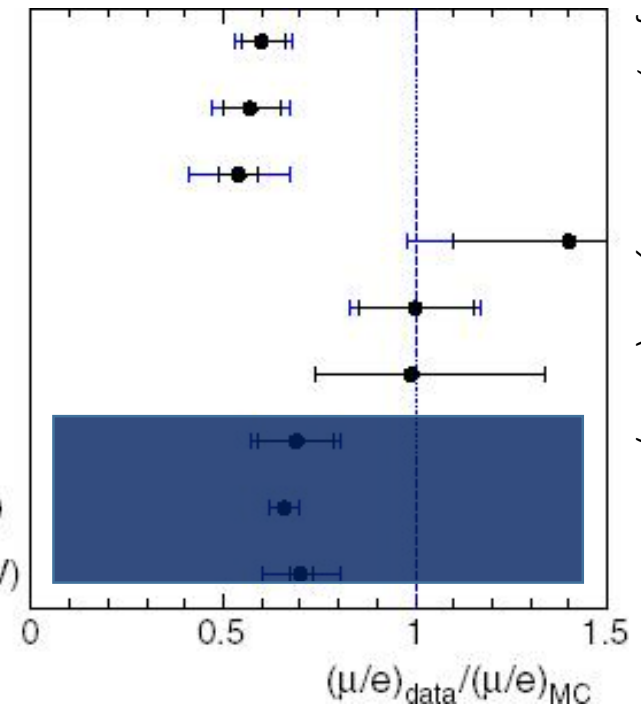


Early measurement of atmospheric ν 's (<1990)

- Tracking calorimeter: Frejus, Nusex, Soudan (hard to find pictures on the web!)
- Water Cherenkov: IMB, Kamiokande
- Measured the number of neutrino interaction in the detector, separating tracks ($=\nu_\mu$) from showers ($=\nu_e$)

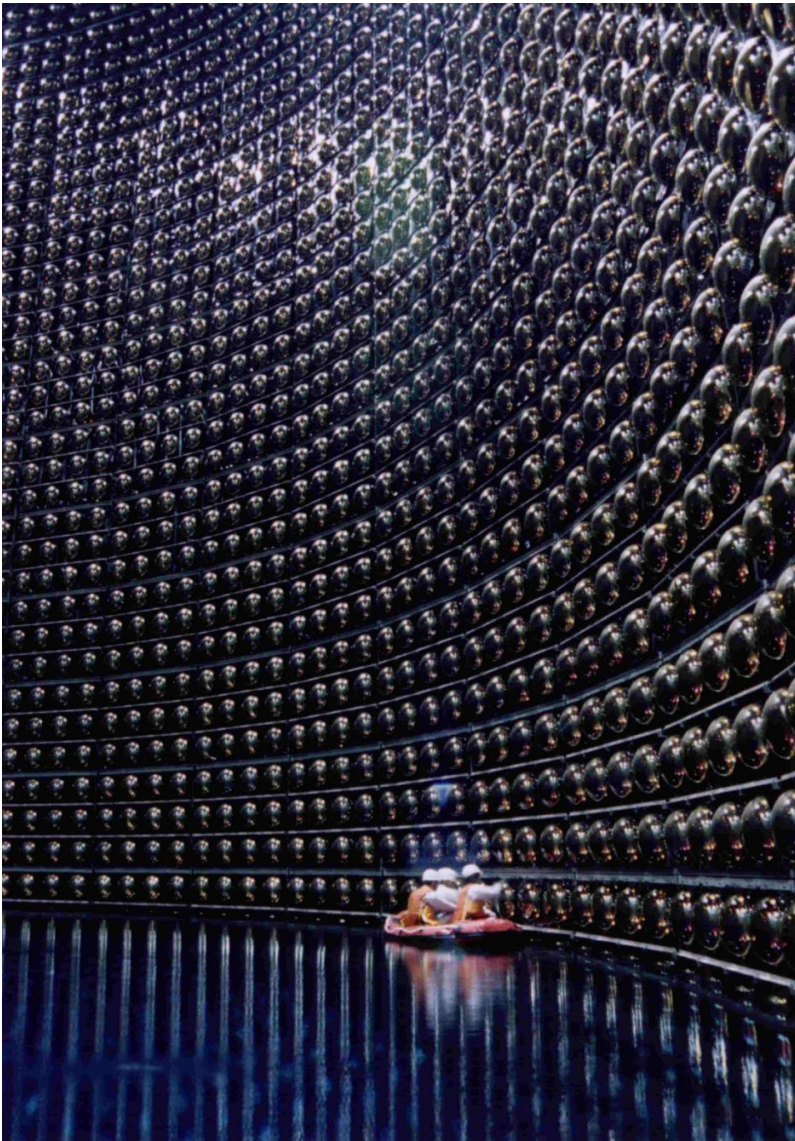


Kam.(sub-GeV)
 Kam.(multi-GeV)
 IMB-3(sub-GeV)
 IMB-3(multi-GeV)
 Frejus
 Nusex
 Soudan-2
 Super-K(sub-GeV)
 Super-K(multi-GeV)



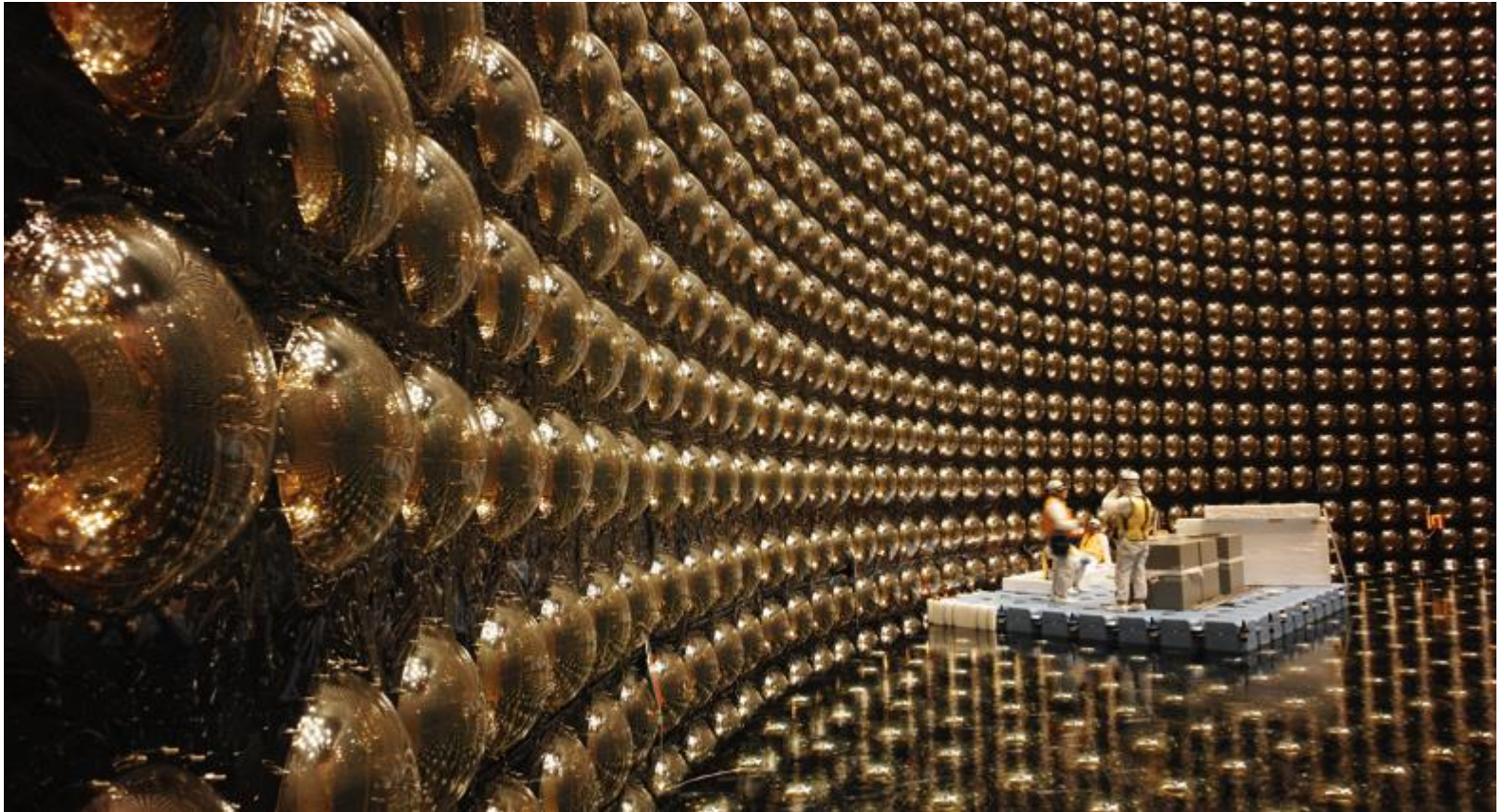
T. Kajita, *New J. Phys.* **6** (2004) 194.

The golden age (1995-2002): SK and MACRO



- SuperKamiokande (SK) is located in Japan, 1000 m Underground (1996 →)
- MACRO was located at LNGS in Italy (1990-2000)

Super-Kamiokande



- SuperKamiokande (SK) is located in Japan, 1000 m Underground
- Active since 1996
- Filled 50.000 ton water
- 11000 large PMTs +2000 PMTs

25 years of Super-Kamiokande

SK-I

1996.4 Start data taking

1998 Evidence of atmospheric ν oscillation (SK)

1999.6 K2K started

2001 Evidence of solar ν oscillation (SNO+SK)

2001.7 data taking was stopped
for detector upgrade

2001.11 Accident

2002.10 data taking was resumed

SK-II

2005 Confirm ν oscillation by accelerator ν (K2K)

SK-III

2006.7 data taking was resumed

2005.10 data taking stopped for full reconstruction

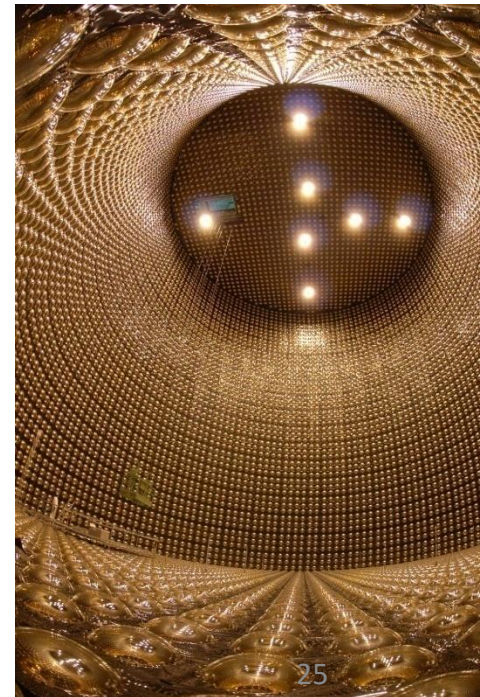
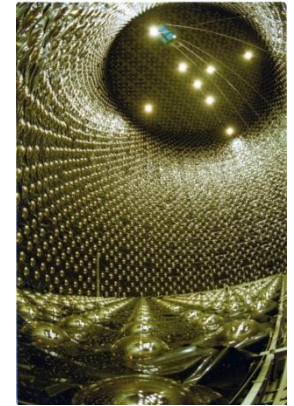
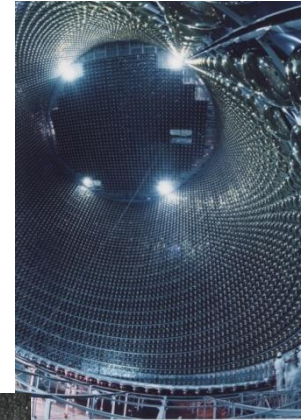
SK-IV

2009 data taking

2018.5 Stop SK-IV

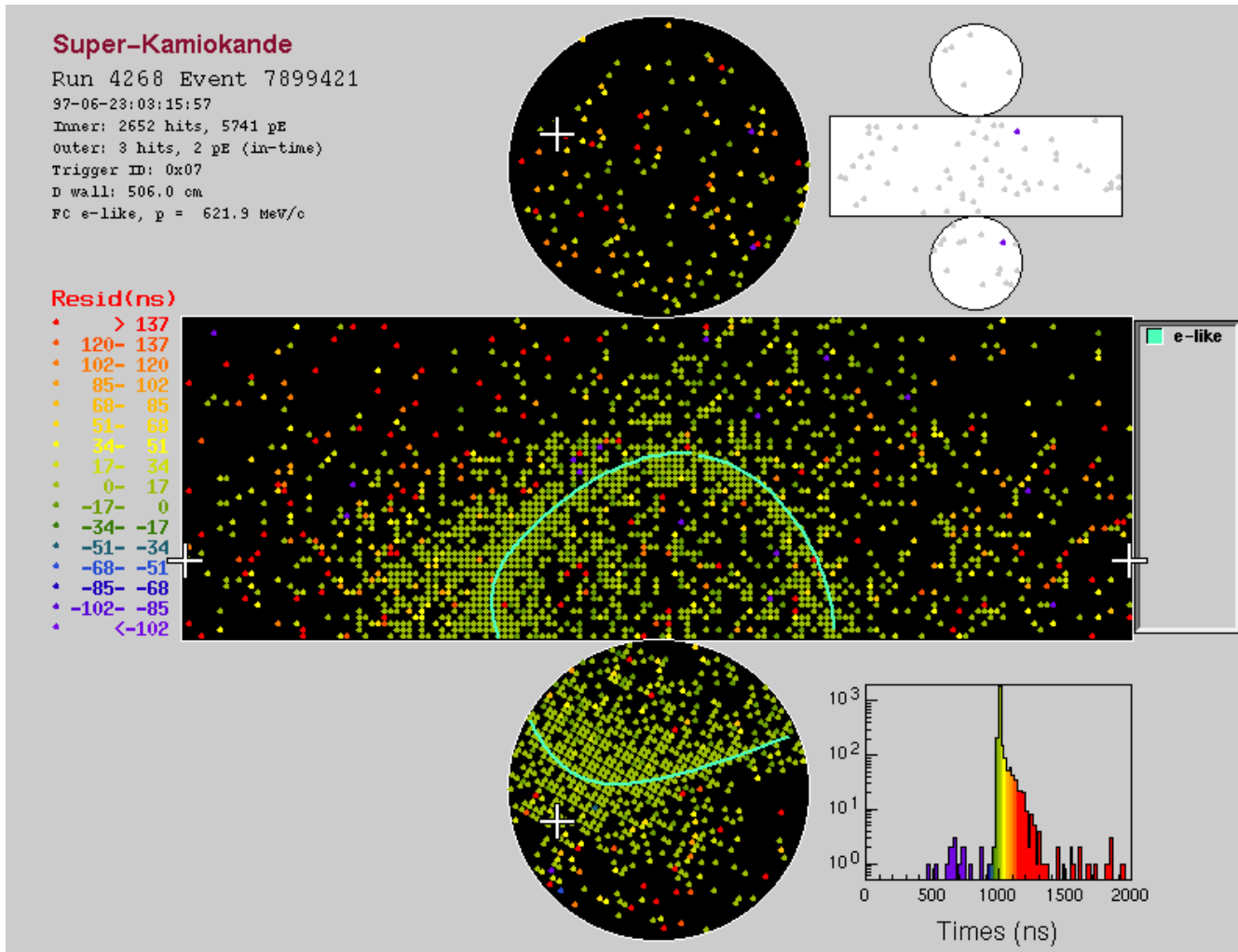
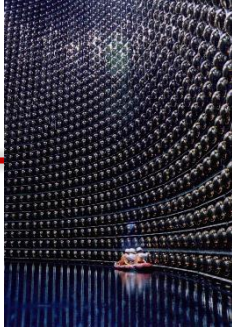
SK-V

2019.1 Start

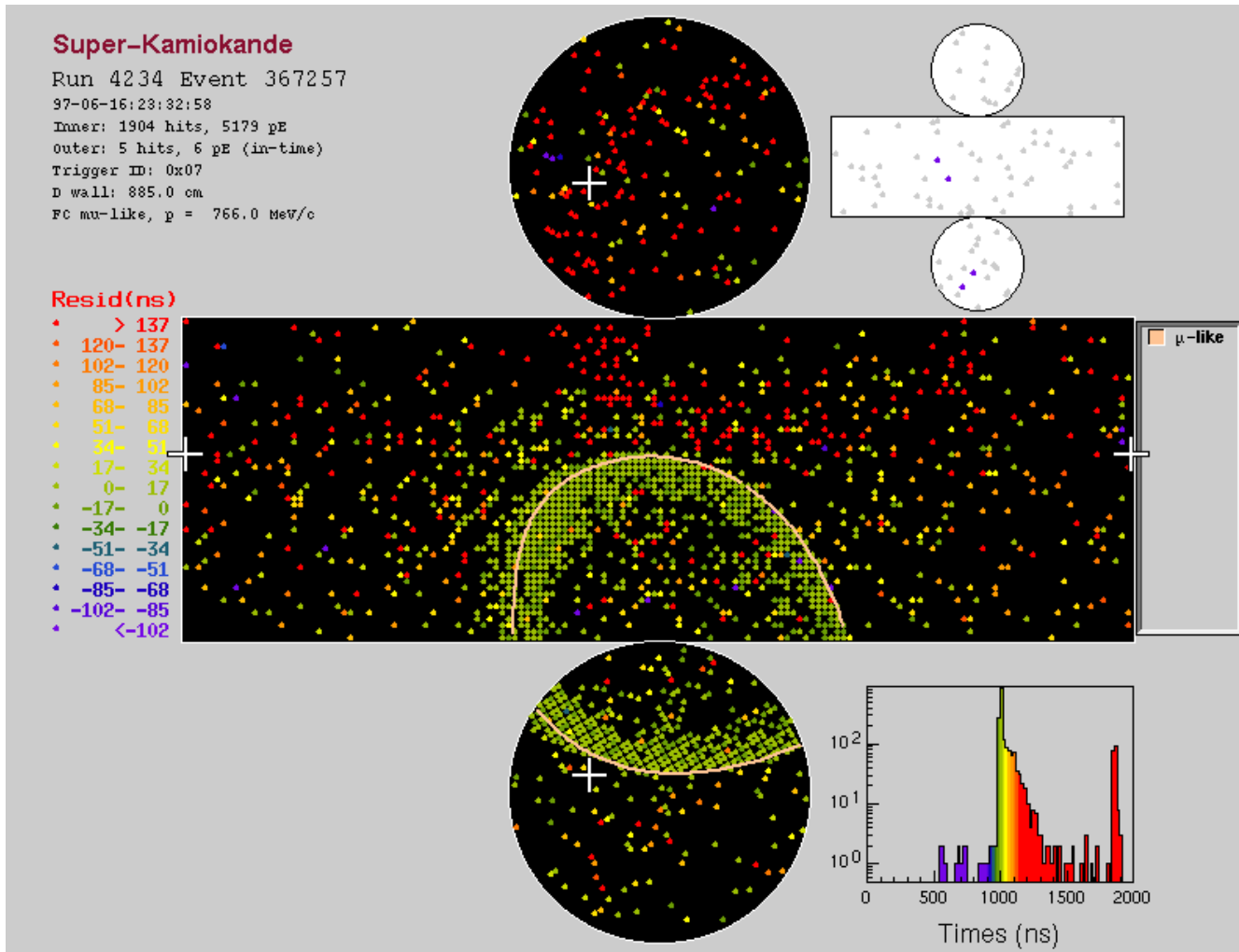
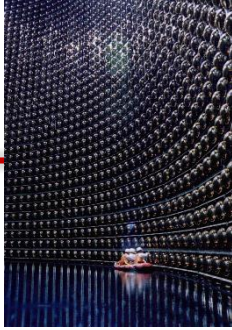


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SuperKamiokande: ν_e



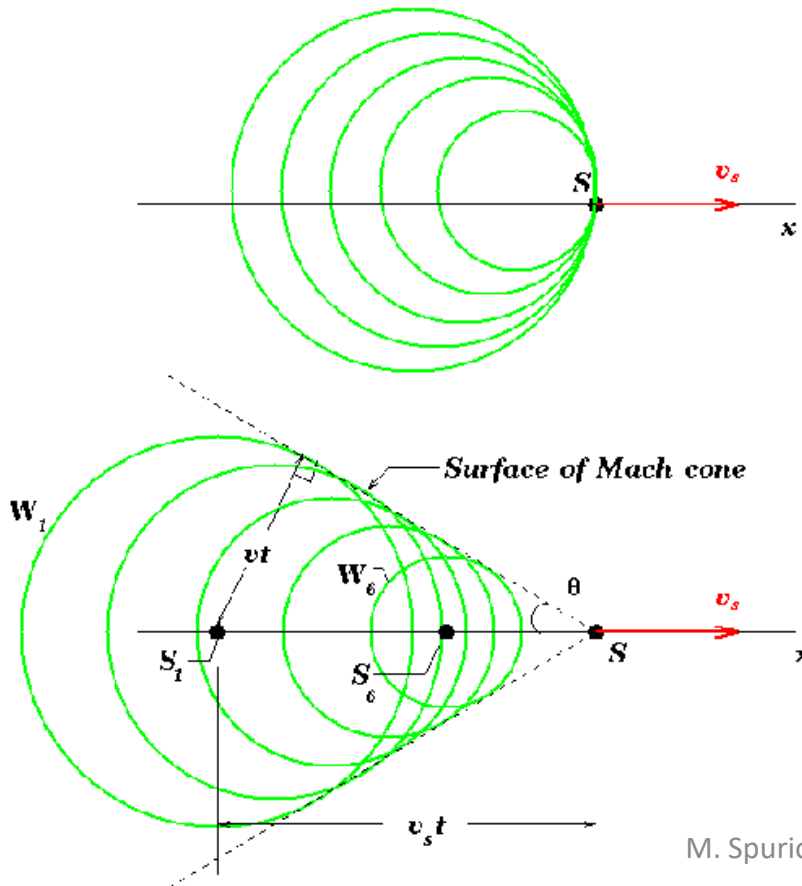
SuperKamiokande: ν_μ



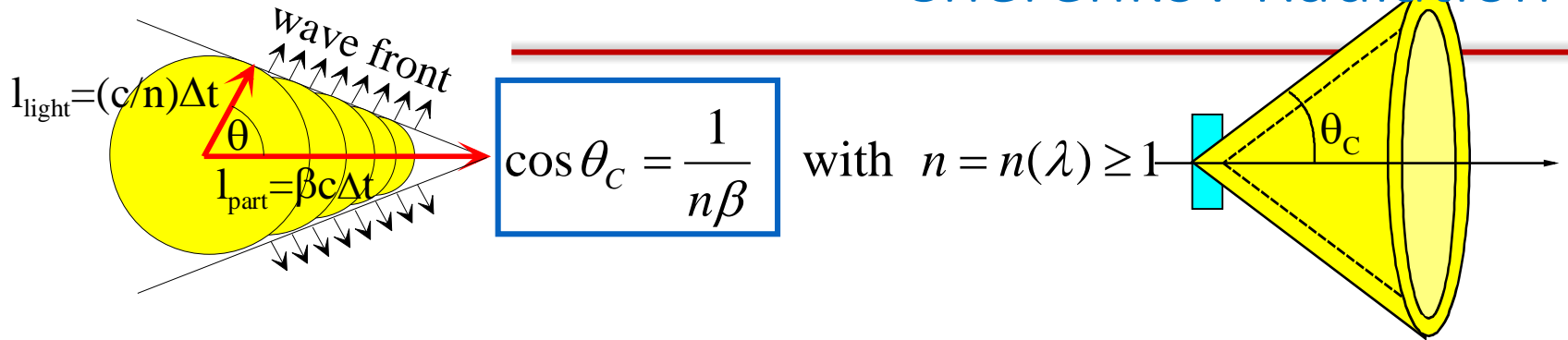
Cherenkov Radiation



- As a charged particle travels, it disrupts the local electromagnetic field (EM) in a medium.
- Electrons in the atoms of the medium will be displaced and polarized by the passing EM field of a charged particle.
- Photons are emitted as an insulator's electrons restore themselves to equilibrium after the disruption has passed.
- In a conductor, the EM disruption can be restored without emitting a photon.
- In normal circumstances, these photons destructively interfere with each other and no radiation is detected.
- However, when the disruption travels faster than light is propagating through the medium, the photons constructively interfere and intensify the observed Cherenkov radiation.



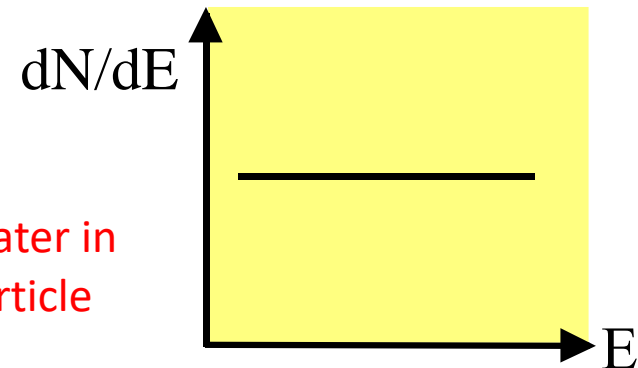
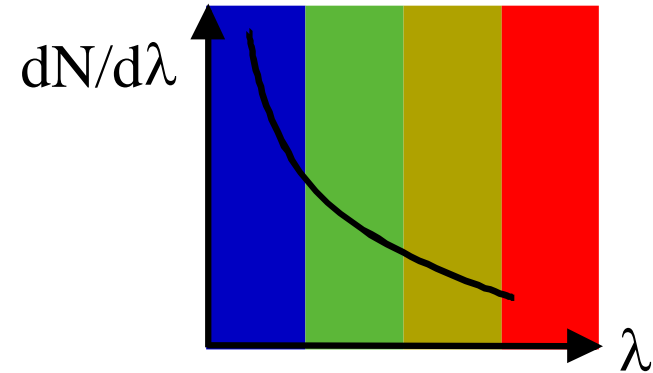
Cherenkov Radiation



- Threshold velocity $\beta_T = 1/n \rightarrow \theta_T \sim 0$
- Angle of emission ($\beta=1$): $\theta_{\text{max}} = \arccos(1/n)$
- Distribution of emitted photons:

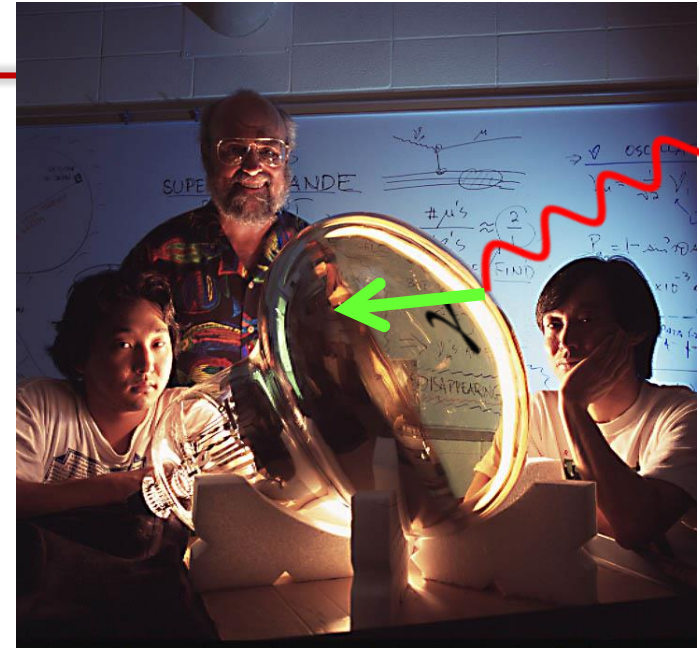
$$\frac{d^2 N}{dx d\lambda} = \frac{2\pi z^2 \alpha}{\lambda^2} \left(1 - \frac{1}{\beta^2 n^2} \right) = \frac{2\pi \cdot z^2 \alpha}{\lambda^2} \sin^2 \theta_c$$

$$\frac{d^2 N}{dx dE} = \frac{z^2 \alpha}{\hbar c} \sin^2 \theta_c$$

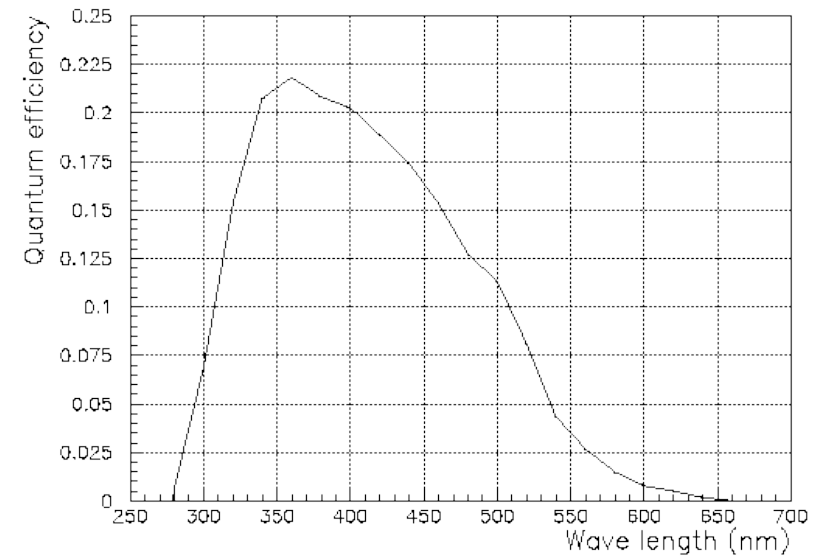


Question 5: Evaluate the number of Cherenkov photons in water in the $\lambda=300-600$ nm interval for a relativistic single charged particle

The PMTs

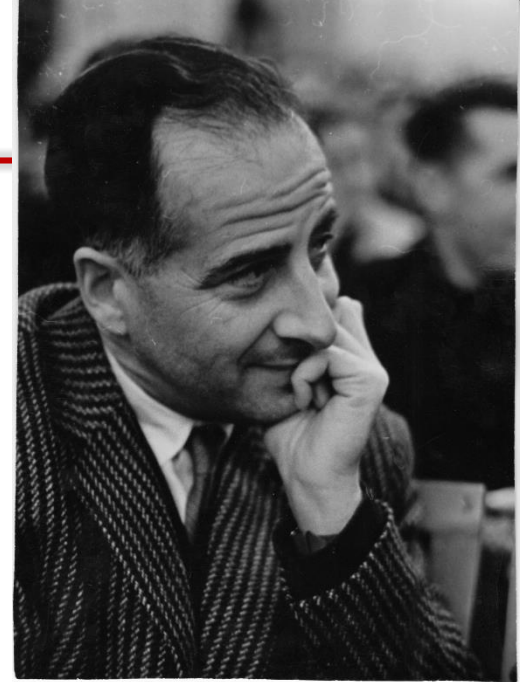


- **Quantum efficiency (Q.e.)**= probability that an incoming photon of a given wavelength produces a signal



The discovery of neutrino masses

- **Neutrino are massless in the Standard Model**
- **Lepton flavour is strictly conserved in SM**
- Idea of neutrinos being massive was first suggested by B. Pontecorvo, from proposal of neutrino oscillations



$|\nu_e\rangle$, $|\nu_\mu\rangle$, $|\nu_\tau\rangle$ =Weak Interactions (WI) eigenstats

$|\nu^1\rangle$, $|\nu^2\rangle$, $|\nu^3\rangle$ =Mass (Hamiltonian) eigenstats

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = V_{\text{PNMS}} \begin{pmatrix} \nu^1 \\ \nu^2 \\ \nu^3 \end{pmatrix}$$

The Pontecorvo-Maki-Nakagawa-Sakata matrix

- We are here

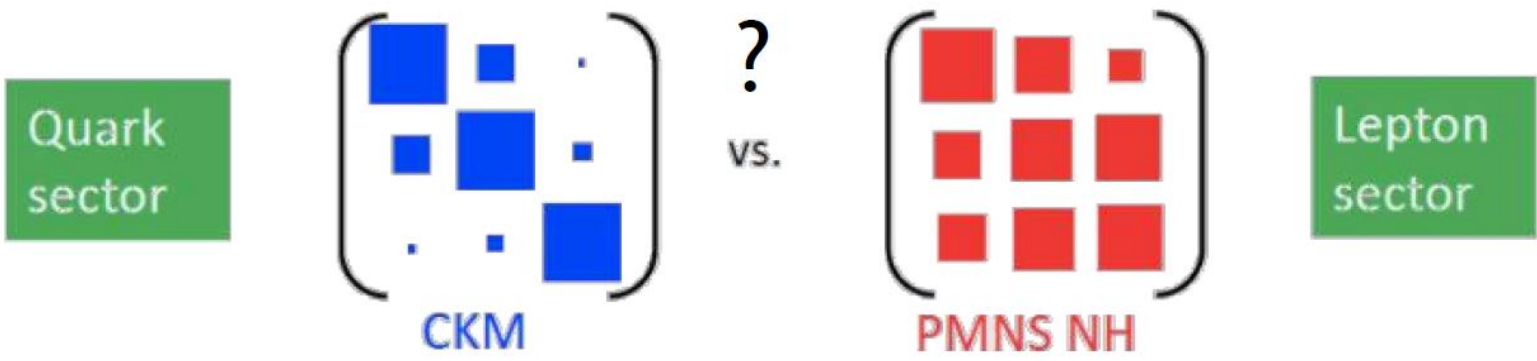
$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13}e^{-i\delta_{CP}} \\ 0 & 1 & 0 \\ -s_{13}e^{i\delta_{CP}} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} e^{i\eta_1} & 0 & 0 \\ 0 & e^{i\eta_2} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

Atmospheric and accelerator $\theta_{23} \approx 50^\circ$
 $|\Delta m_{32}^2| \sim 2.5 \times 10^{-3} \text{eV}^2$

Reactor and accelerator $\theta_{13} \approx 8^\circ$
 Accelerator only: $\delta_{CP} = ?$

Solar and reactor $\theta_{12} \approx 34^\circ$
 $\Delta m_{12}^2 \sim 7.5 \times 10^{-5} \text{eV}^2$

Normal hierarchy Inverted hierarchy

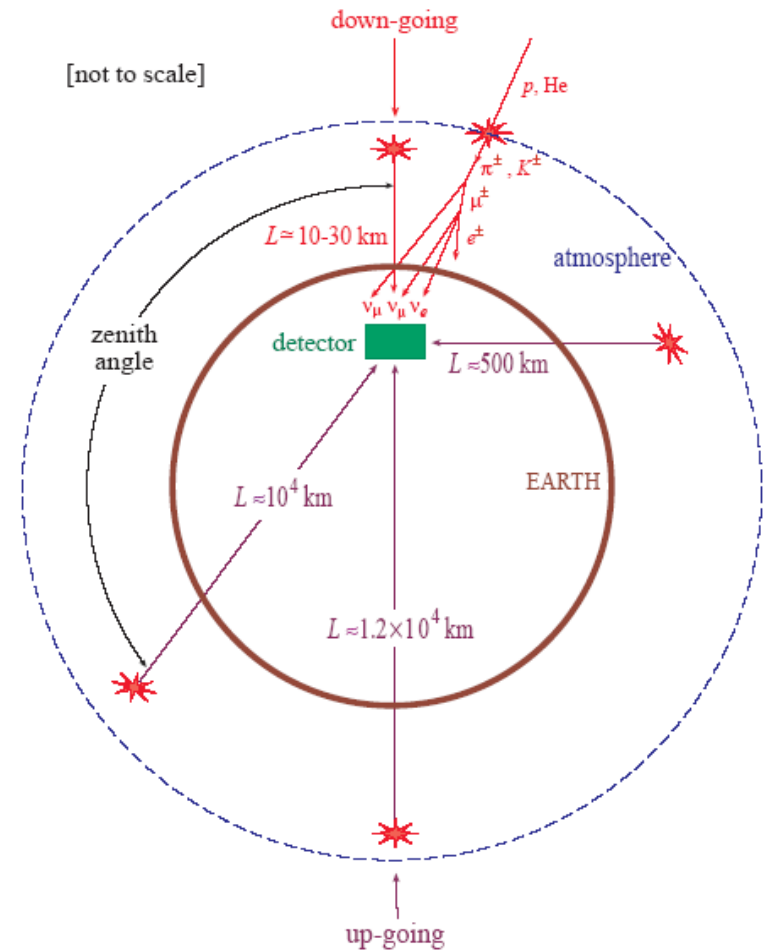


...with atmospheric neutrinos

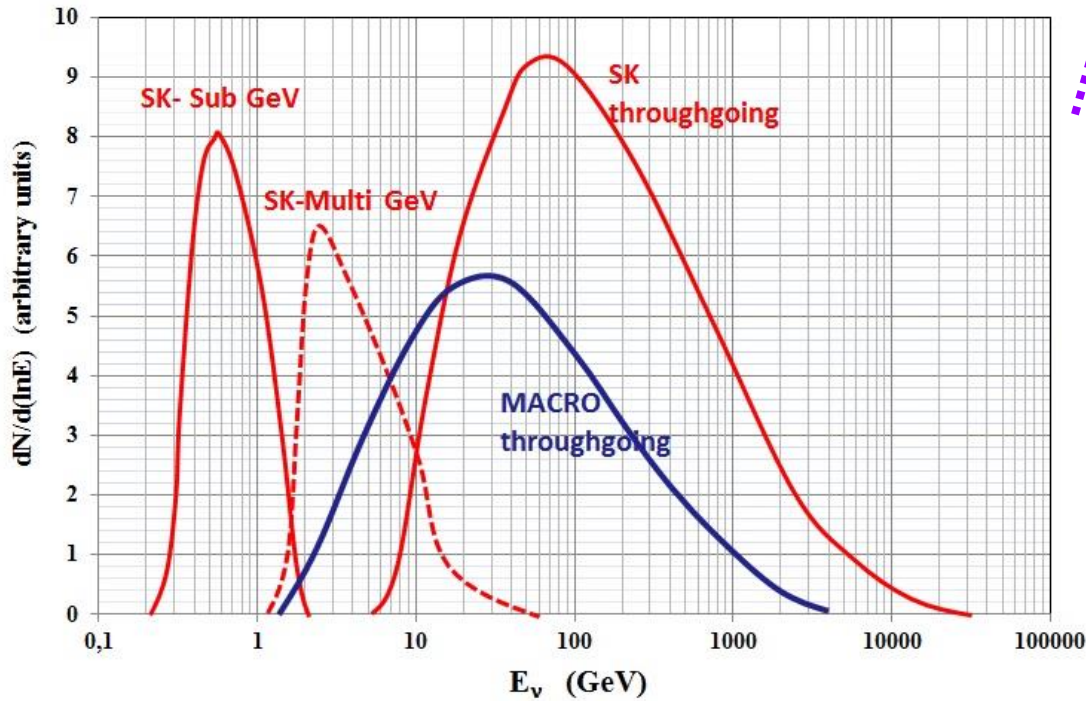
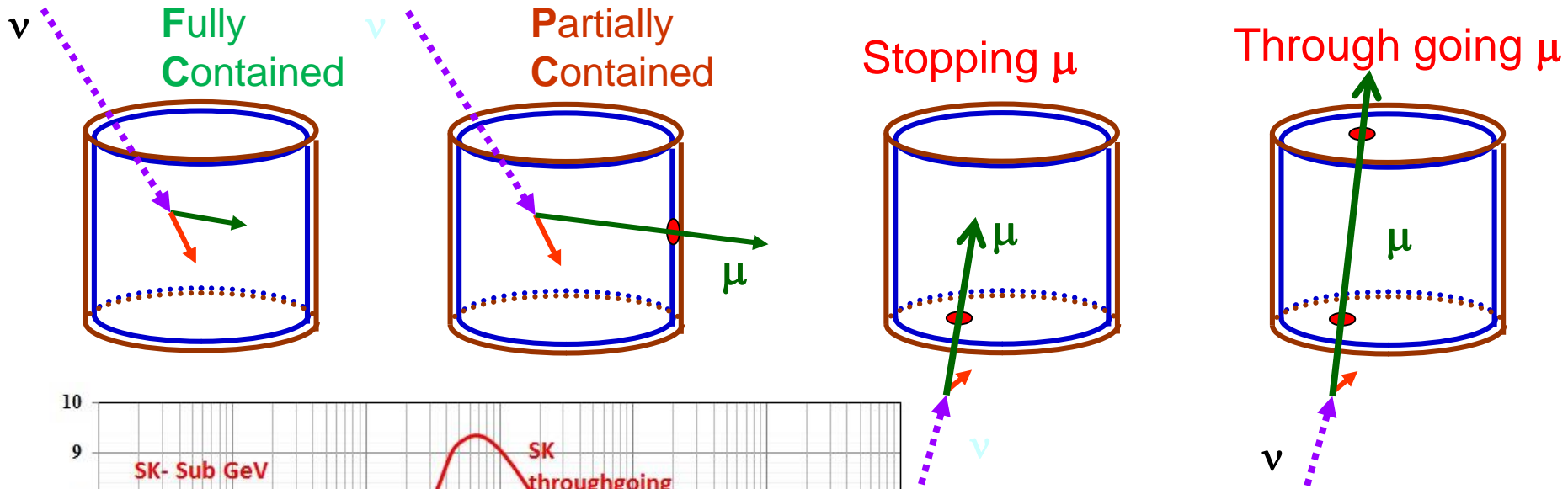


$$P_{\nu_{\mu} \rightarrow \nu_{\mu}} = 1 - \sin^2 2\vartheta_{23} \sin^2 \left[1.27 \frac{\Delta m_{23}^2 L}{E_{\nu}} \right]$$

- $\Delta m_{23}^2, \vartheta_{23} \rightarrow$ from Nature;
- E_{ν} = experimental parameter (energy distribution of neutrino giving a particular configuration of events)
- L = experimental parameter (neutrino path length from production to interaction)



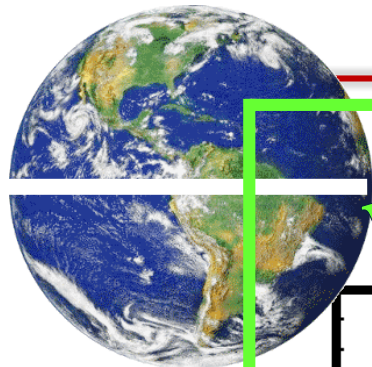
ν energy/direction: event topology



Energy spectrum (Monte Carlo) of atmospheric ν seen with different event topologies (Super-Kamiokande, MACRO)

$\cos\Theta > 0$

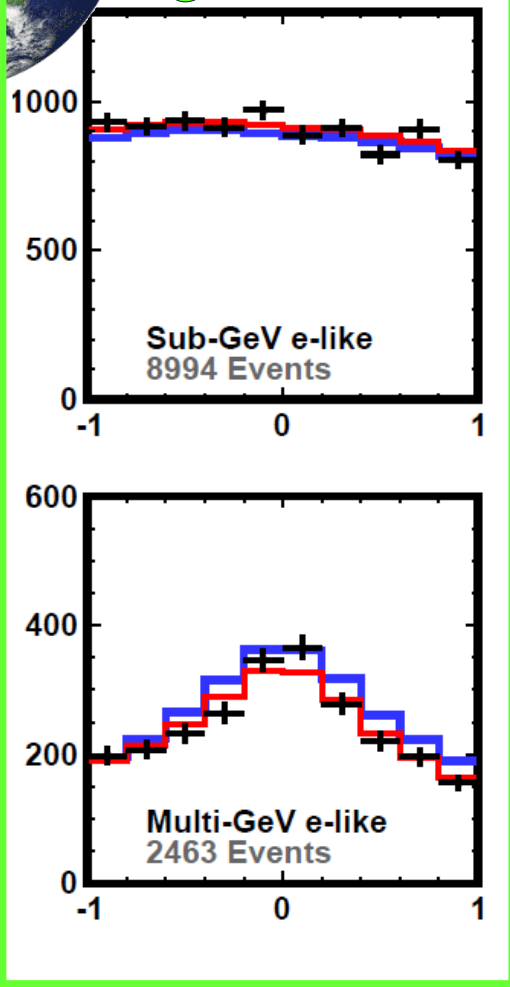
SuperKamiokande I-IV: results



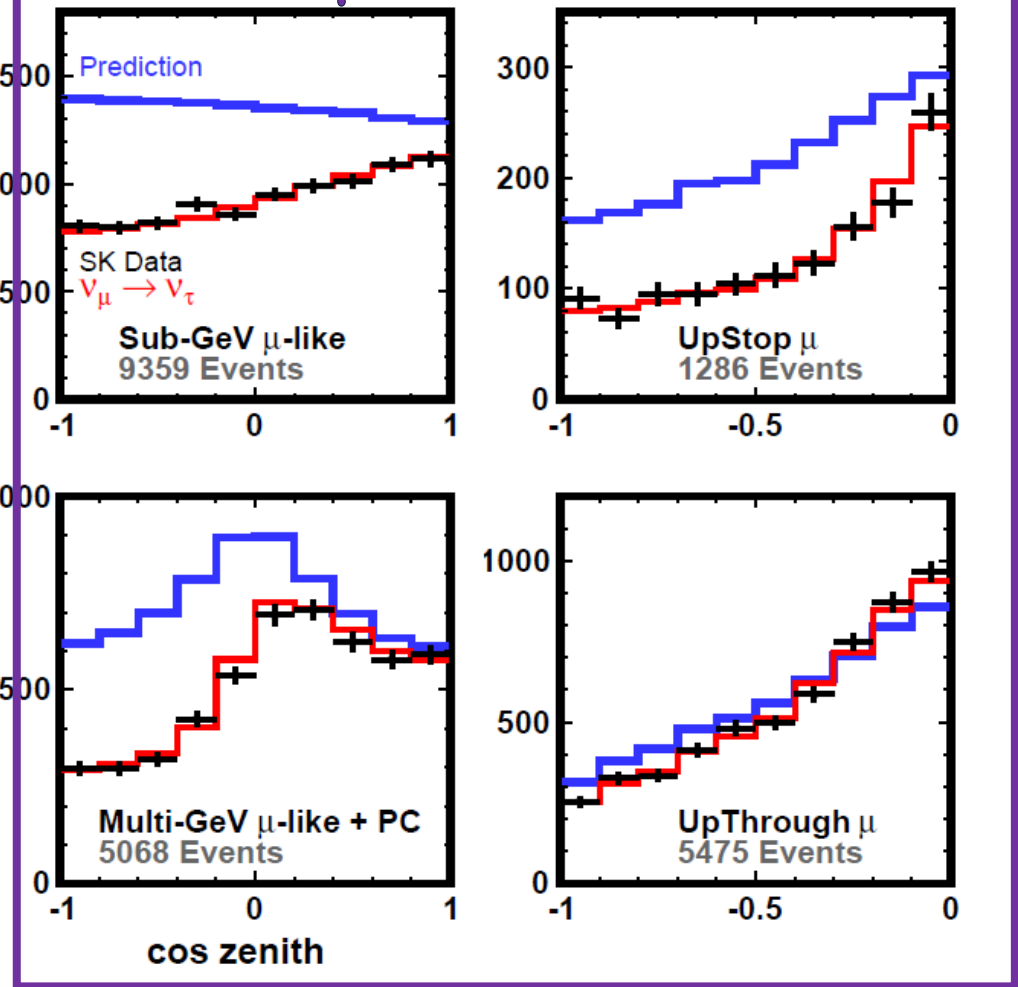
$\cos\Theta < 0$

ν_e LIKE

Number of Events

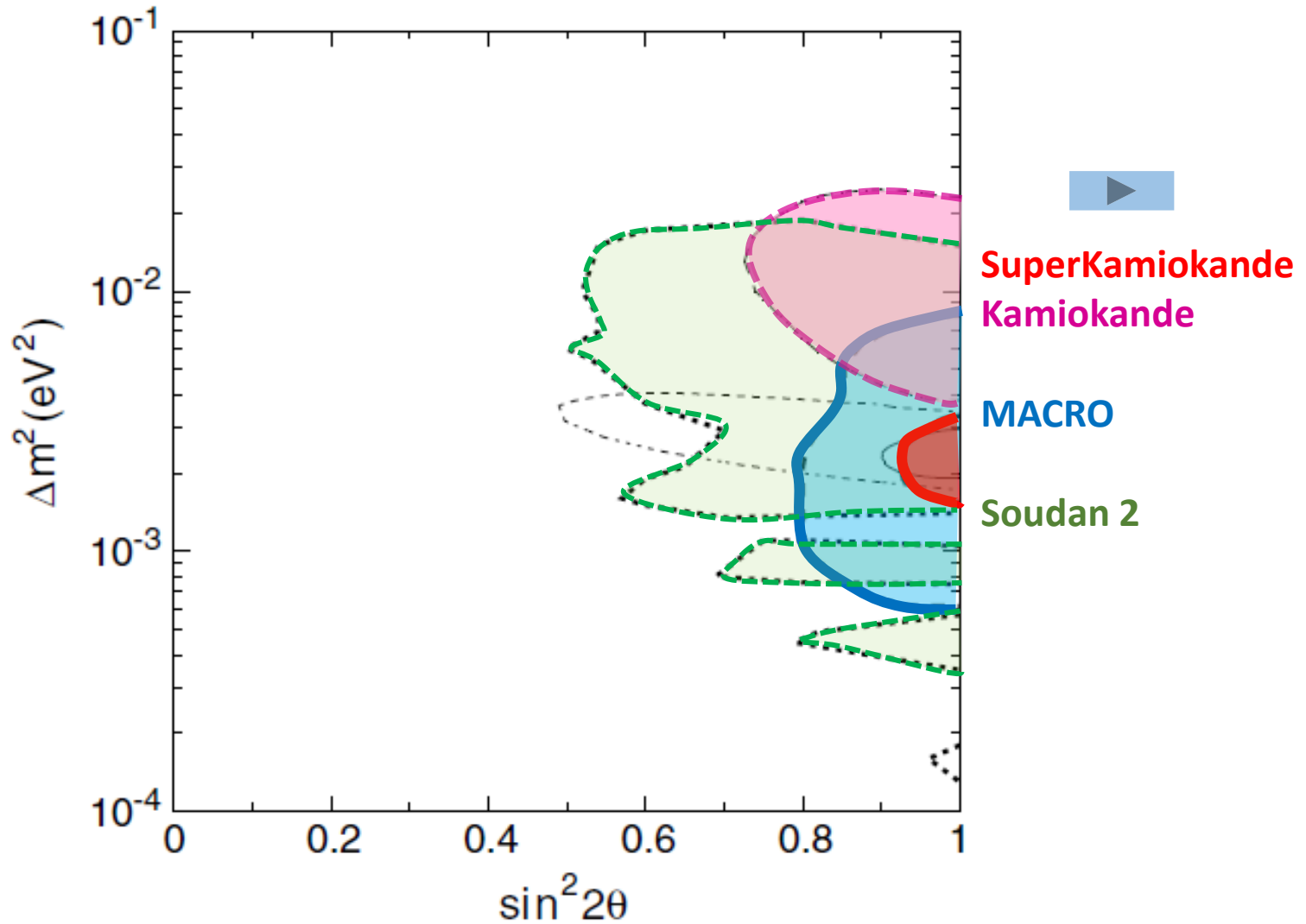


ν_μ LIKE

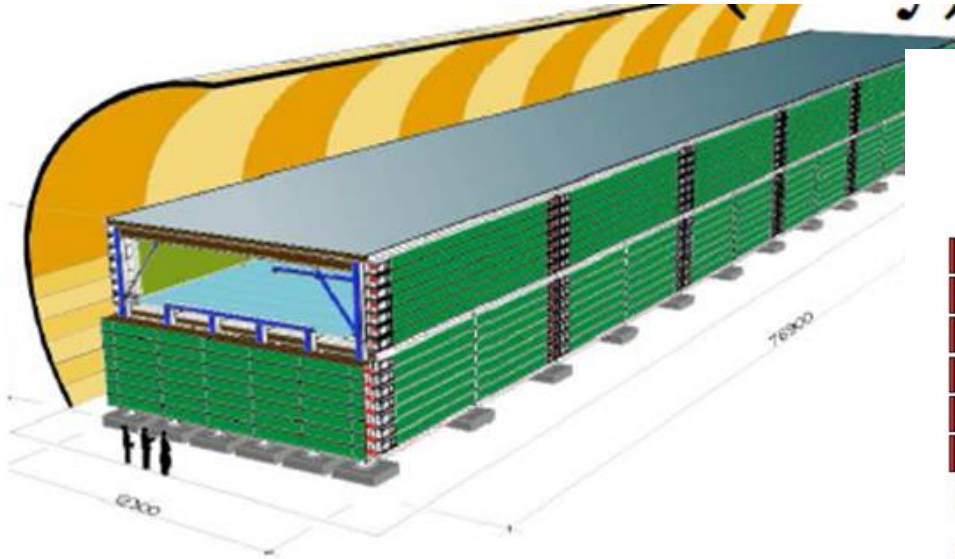


arXiv:1412.5234v1

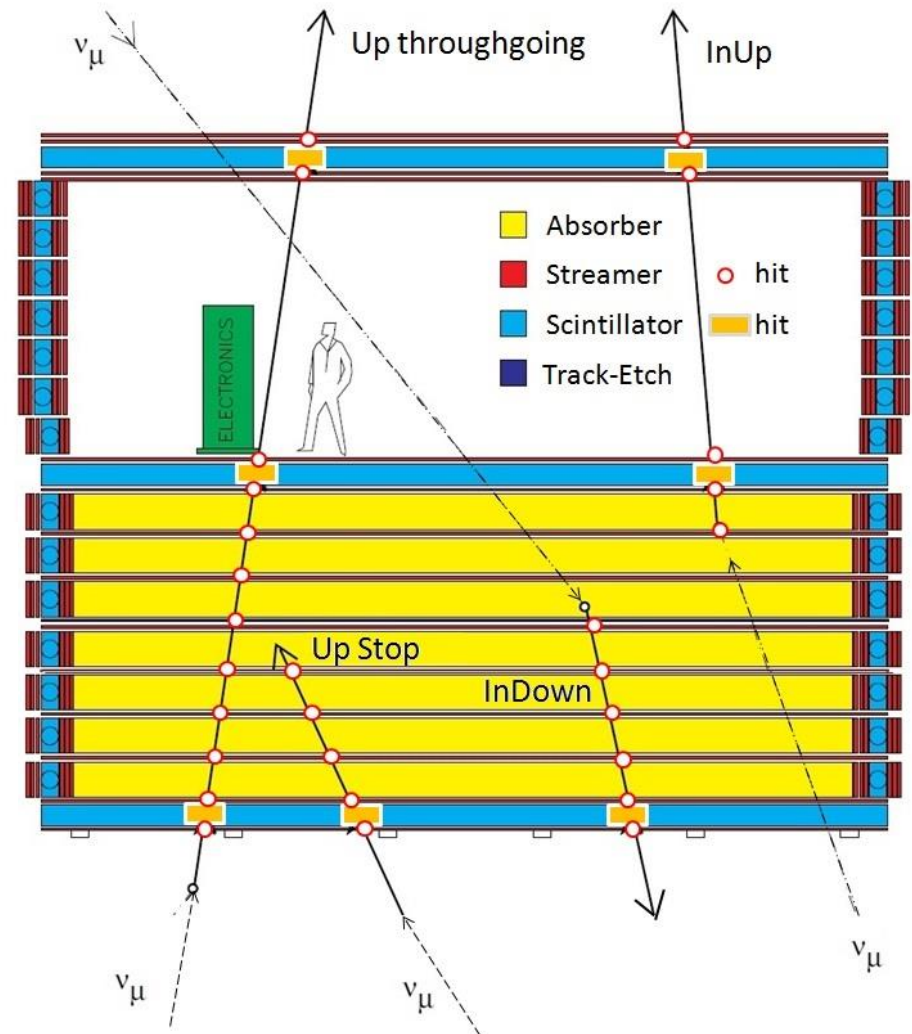
Discovery of neutrino oscillations



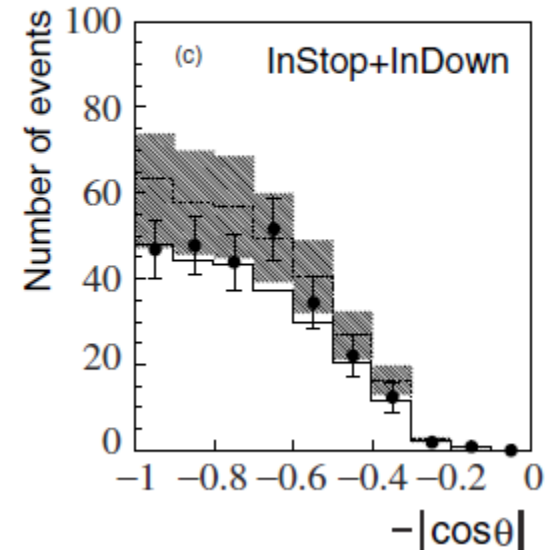
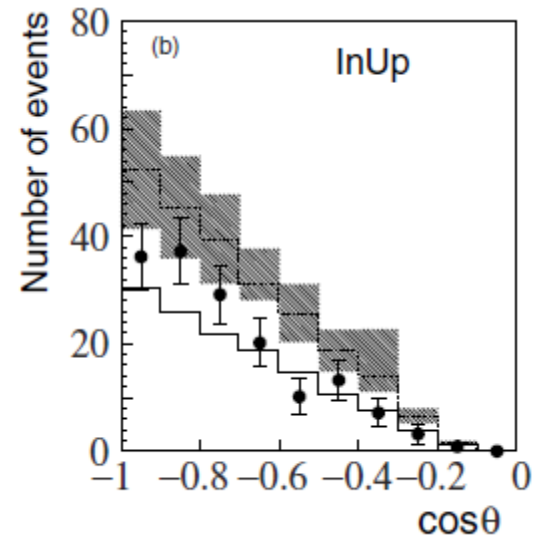
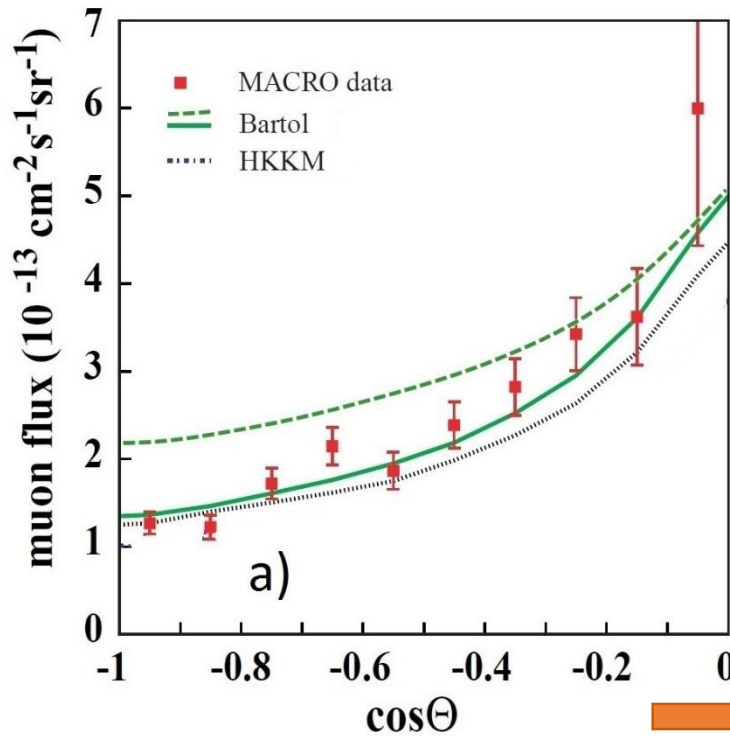
MACRO@ LNGS



- **Liquid scintillator counters**, (3 planes) for the measurement of **time** and **dE/dx**.
- **Streamer tubes** (14 planes), for the measurement of the track position ($<1^\circ$);
- Detector mass: 5.3 kton
- 1 Up going muons $\sim 10^6$ downgoing muons
- Different neutrino topologies

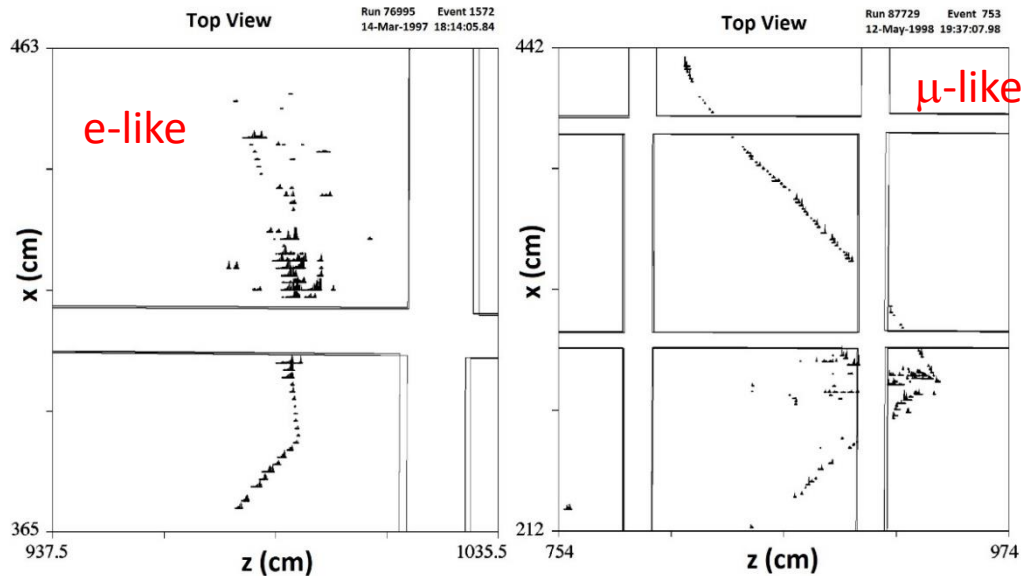


The MACRO neutrino deficits

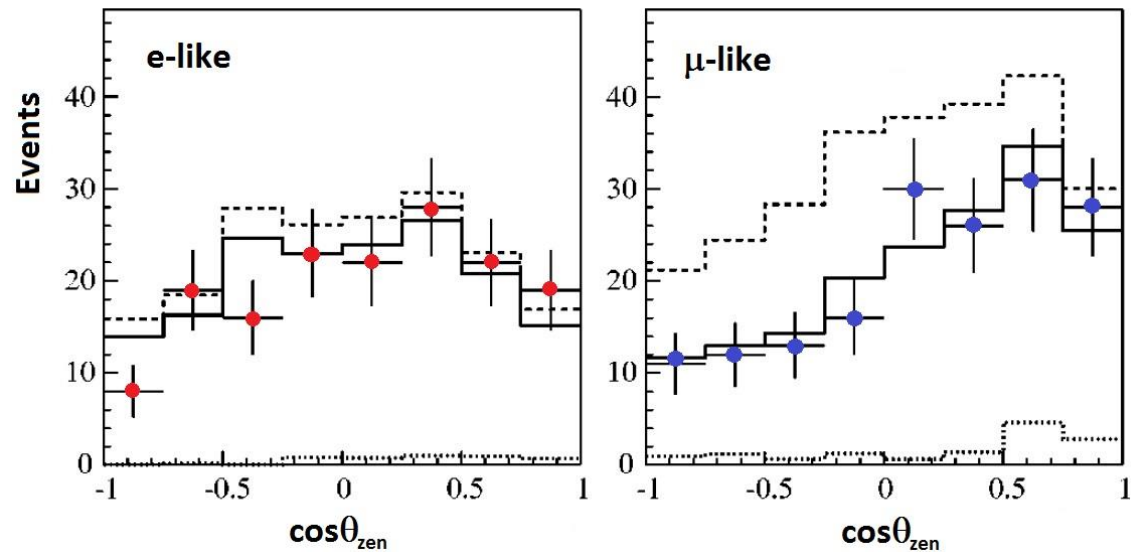


- Completely different topology w.r.t. SK
- Different experimental technique
- Deformation of the angular distribution w.r.t. expectation
- Missing events from the vertical direction
- Interpretation: oscillations
- The same oscillation parameters!

The Soudan II neutrino deficit

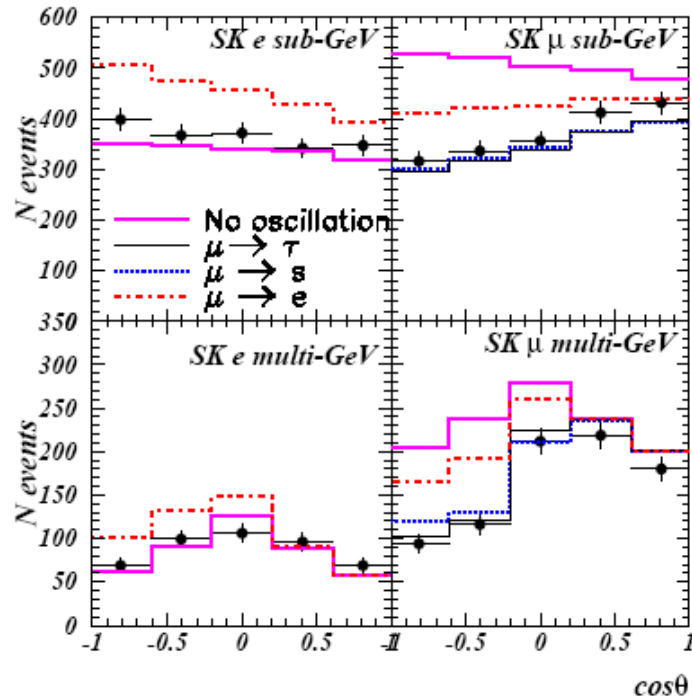


- Iron tracking calorimeter
- 770 t fiducial mass
- Active from 1989 to 2001 in the Soudan Mine (USA)
- (P)contained events
- μ -like deficit from below

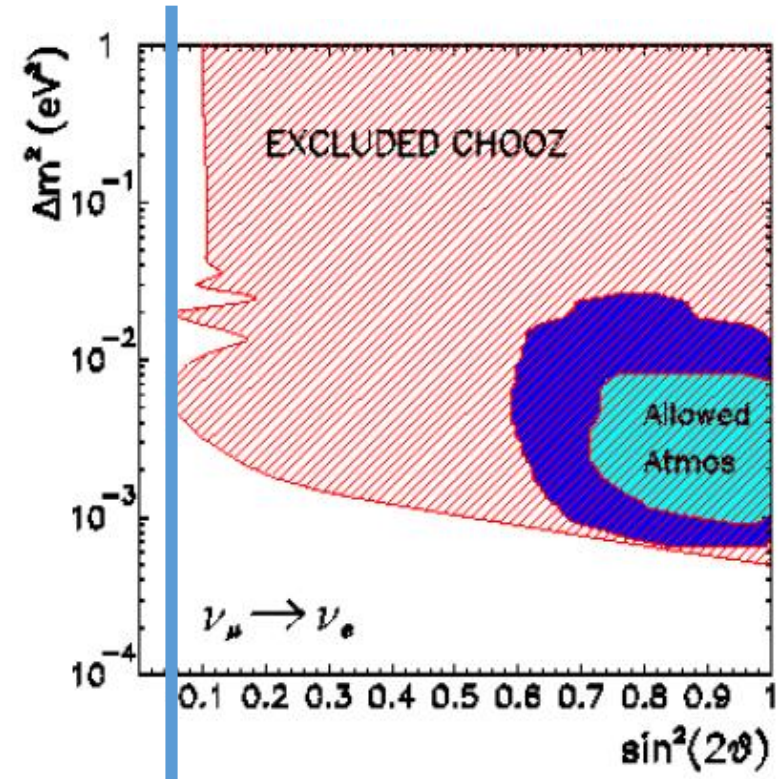


Why not $\nu_\mu \rightarrow \nu_e$?

- i.e., through the ϑ_{13} angle of the PMNS matrix



Apollonio et al., CHOOZ Coll.,
 Phys.Lett.B466 (1999) 415



PDG value:
 0.095 ± 0.010

The 2015 Nobel Prize

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2015

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► Takaaki Kajita

► Arthur B. McDonald

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All Nobel Prizes in 2015



The Nobel Prize in Physics 2015

Takaaki Kajita, Arthur B. McDonald

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The Nobel Prize in Physics 2015



Photo © Takaaki Kajita

Takaaki Kajita

Prize share: 1/2

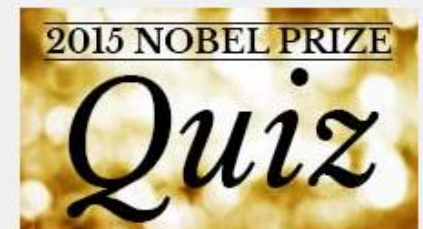


Photo: K. McFarlane,
Queen's University
/SNOLAB

Arthur B. McDonald

Prize share: 1/2

The Nobel Prize in Physics 2015 was awarded jointly to Takaaki Kajita and Arthur B. McDonald *"for the discovery of neutrino oscillations, which shows that neutrinos have mass"*



Evidence for oscillation of atmospheric neutrinos

The Super-Kamiokande Collaboration

Y.Fukuda^a, T.Hayakawa^a, E.Ichihara^a, K.Inoue^a, K.Ishihara^a, H.Ishino^a, Y.Itow^a, T.Kajita^a,
J.Kameda^a, S.Kasuga^a, K.Kobayashi^a, Y.Kobayashi^a, Y.Koshio^a, M.Miura^a, M.Nakahata^a,
S.Nakayama^a, A.Okada^a, K.Okumura^a, N.Sakurai^a, M.Shiozawa^a, Y.Suzuki^a, Y.Takeuchi^a, Y.Totsuka^a,
S.Yamada^a, M.Earl^b, A.Habig^b, E.Kearns^b, M.D.Messier^b, K.Scholberg^b, J.L.Stone^b, L.R.Sulak^b,
C.W.Walter^b, M.Goldhaber^c, T.Barszczak^d, D.Casper^d, W.Gajewski^d, P.G.Halverson^{d,*}, J.Hsu^d,
W.R.Kropp^d, L.R. Price^d, F.Reines^d, M.Smy^d, H.W.Sobel^d, M.R.Vagins^d, K.S.Ganezer^e, W.E.Keig^e,
R.W.Ellsworth^f, S.Tasaka^g, J.W.Flanagan^{h,†}, A.Kibayashi^h, J.G.Learned^h, S.Matsuno^h, V.J.Stenger^h,
D.Takemori^h, T.Ishiiⁱ, J.Kanzakiⁱ, T.Kobayashiⁱ, S.Mineⁱ, K.Nakamuraⁱ, K.Nishikawaⁱ, Y.Oyamaⁱ,
A.Sakaiⁱ, M.Sakudaⁱ, O.Sasakiⁱ, S.Echigo^j, M.Kohama^j, A.T.Suzuki^j, T.J.Haines^{k,d}, E.Blaufuss^l,
B.K.Kim^l, R.Sanford^l, R.Svoboda^l, M.L.Chen^m, Z.Conner^{m,‡}, J.A.Goodman^m, G.W.Sullivan^m, J.Hillⁿ,
C.K.Jungⁿ, K.Martensⁿ, C.Maugerⁿ, C.McGrewⁿ, E.Sharkeyⁿ, B.Virenⁿ, C.Yanagisawaⁿ, W.Doki^o,
K.Miyano^o, H.Okazawa^o, C.Saji^o, M.Takahata^o, Y.Nagashima^p, M.Takita^p, T.Yamaguchi^p,
M.Yoshida^p, S.B.Kim^q, M.Etoh^r, K.Fujita^r, A.Hasegawa^r, T.Hasegawa^r, S.Hatakeyama^r, T.Iwamoto^r,
M.Koga^r, T.Maruyama^r, H.Ogawa^r, J.Shirai^r, A.Suzuki^r, F.Tsushima^r, M.Koshihara^s, M.Nemoto^t,
K.Nishijima^t, T.Futagami^u, Y.Hayato^{u,§}, Y.Kanaya^u, K.Kaneyuki^u, Y.Watanabe^u, D.Kielczewska^{v,d},
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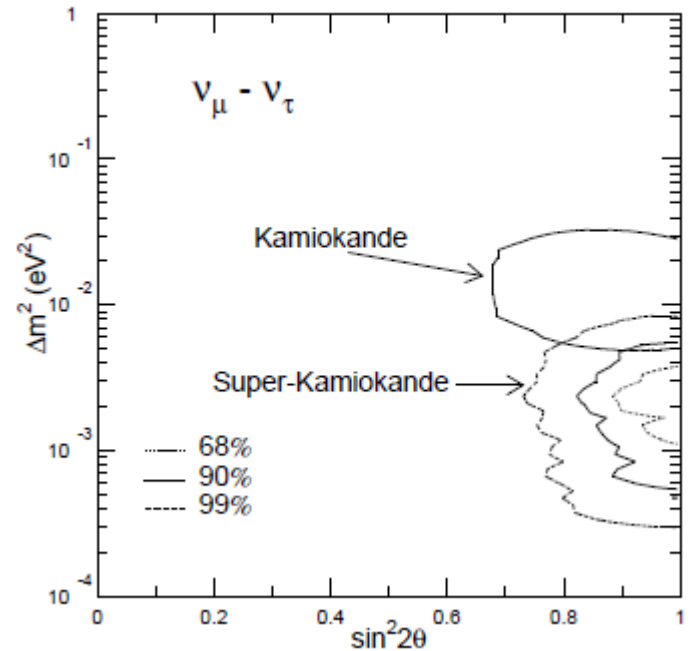
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^uDepartment of Physics, Tokyo Institute of Technology, Meguro, Tokyo 152-

^vInstitute of Experimental Physics, Warsaw University, 00-681 Warsaw,

^wDepartment of Physics, University of Washington, Seattle, WA 98195-15

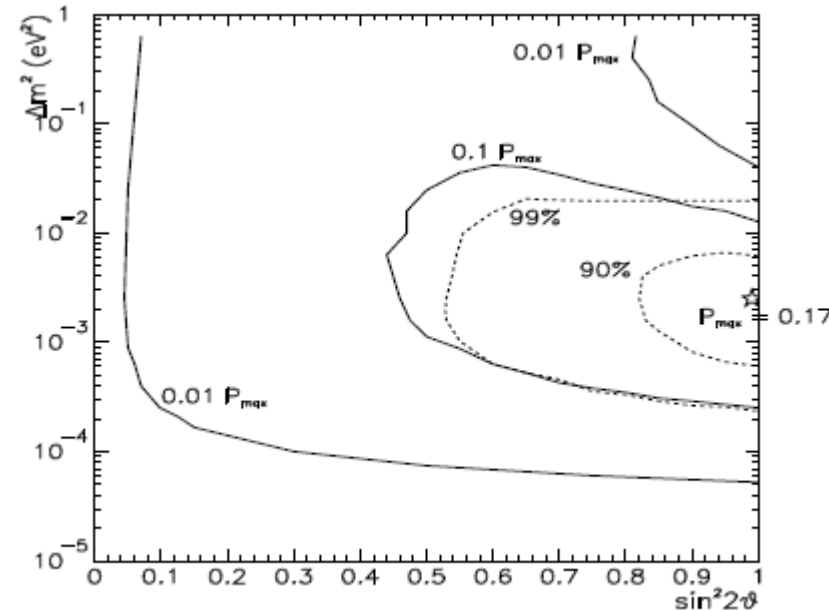


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M. Ambrosio¹², R. Antolini⁷, C. Aramo^{7,n}, G. Auriemma^{14,a}, A. Baldini¹³, G. C. Barbarino¹², B. C. Barish⁴, G. Battistoni^{6,b}, R. Bellotti¹, C. Bemporad¹³, P. Bernardini¹⁰, H. Bilokon⁶, V. Bisi¹⁶, C. Bloise⁶, C. Bower⁸, S. Bussino¹⁴, F. Cafagna¹, M. Calicchio¹, D. Campana¹², M. Carboni⁶, M. Castellano¹, S. Cecchini^{2,c}, F. Cei^{11,13}, V. Chiarella⁶, B. C. Choudhary⁴, S. Coutu^{11,o}, L. De Benedictis¹, G. De Cataldo¹, H. Dekhissi^{2,17}, C. De Marzo¹, I. De Mitrì⁹, J. Derkaoui^{2,17}, M. De Vincenzi^{14,e}, A. Di Credico⁷, O. Erriquez¹, C. Favuzzi¹, C. Forti⁶, P. Fusco¹, G. Giacomelli², G. Giannini^{13,f}, N. Giglietto¹, M. Giorgini², M. Grassi¹³, L. Gray^{4,7}, A. Grillo⁷, F. Guarino¹², P. Guarnaccia¹, C. Gustavino⁷, A. Habig³, K. Hanson¹¹, A. Hawthorne⁸, R. Heinz⁸, Y. Huang⁴, E. Iarocci^{6,g}, E. Katsavounidis⁴, I. Katsavounidis⁴, E. Kearns³, H. Kim⁴, S. Kyriazopoulou⁴, E. Lamanna¹⁴, C. Lane⁵, D. S. Levin¹¹, P. Lipari¹⁴, N. P. Longley^{4,l}, M. J. Longo¹¹, F. Maaroufi^{2,17}, G. Mancarella¹⁰, G. Mandrioli², S. Manzoor^{2,m}, A. Margiotta Neri², A. Marini⁶, D. Martello¹⁰, A. Marzari-Chiesa¹⁶, M. N. Mazziotta¹, C. Mazzotta¹⁰, D. G. Michael⁴, S. Mikheyev^{4,7,h}, L. Miller⁸, P. Monacelli⁹, T. Montaruli¹, M. Monteno¹⁶, S. Mufson⁸, J. Musser⁸, D. Nicolò^{13,d}, R. Nolty⁴, C. Okada³, C. Orth³, G. Osteria¹², M. Ouchrif^{2,17}, O. Palamara¹⁰, V. Patera^{6,g}, L. Patrizii², R. Pazzi¹³, C. W. Peck⁴, S. Petrera⁹, P. Pistilli^{14,e}, V. Popa^{2,i}, V. Pugliese¹⁴, A. Rainó¹, J. Reynoldson⁷, F. Ronga⁶, U. Rubizzo¹², A. Sanzgiri¹⁵, C. Satriano^{14,a}, L. Satta^{6,g}, E. Scapparone⁷, K. Scholberg³, A. Sciubba^{6,g}, P. Serra-Lugaresi², M. Severi¹⁴, M. Sioli², M. Sitta¹⁶, P. Spinelli¹, M. Spinetti⁶, M. Spurio², R. Steinberg⁵, J. L. Stone³, L. R. Sulak³, A. Surdo¹⁰, G. Tarlé¹¹, V. Togo², D. Ugolotti², M. Vakili¹⁵, C. W. Walter³, arXiv:hep-ex/9807005v1

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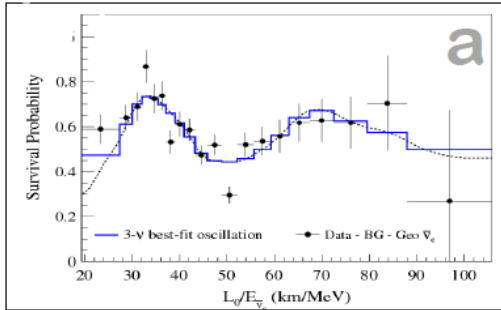
a Also Università della Basilicata, 85100 Potenza, Italy
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 k Also INFN Roma, 00146 Roma, Italy
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 o Also INFN Roma, 00146 Roma, Italy



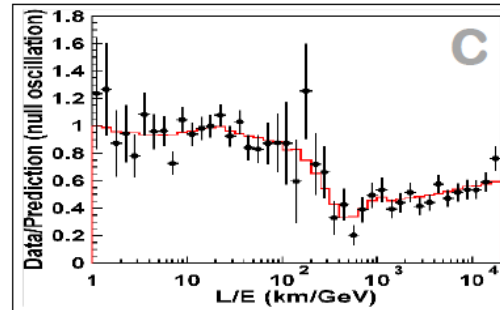
Now and next... [From E. Lisi (INFN-Bari)]



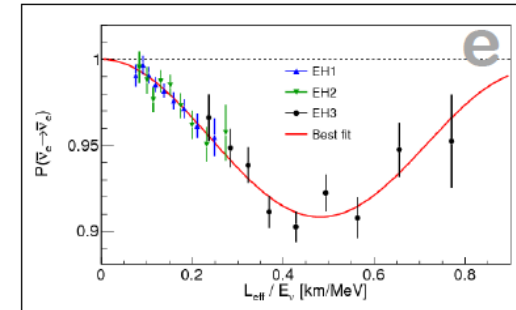
$e \rightarrow e$ (KamLAND, KL)



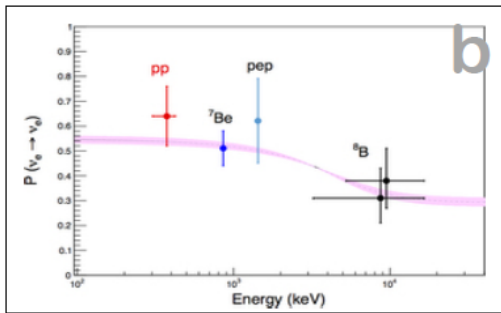
$\mu \rightarrow \mu$ (Atmospheric)



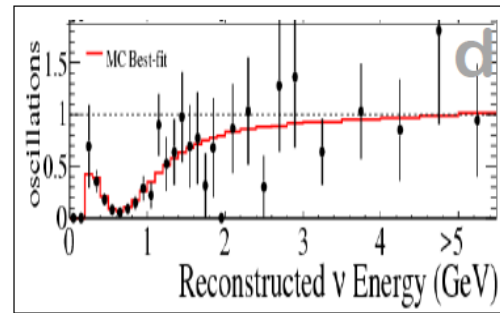
$e \rightarrow e$ (SBL React.)



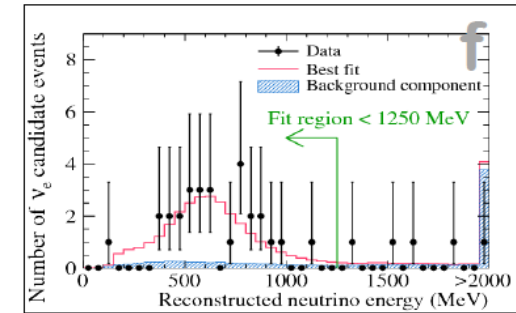
$e \rightarrow e$ (Solar)



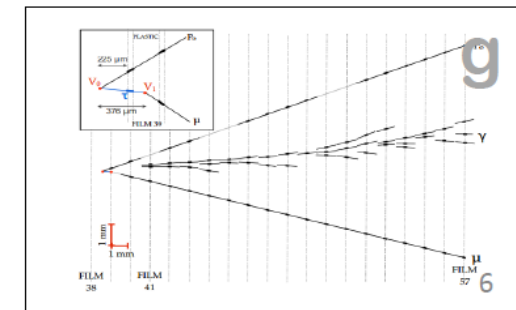
$\mu \rightarrow \mu$ (LBL Accel)



$\mu \rightarrow e$ (LBL Accel)



$\mu \rightarrow \tau$ (OPERA, SK, DC)

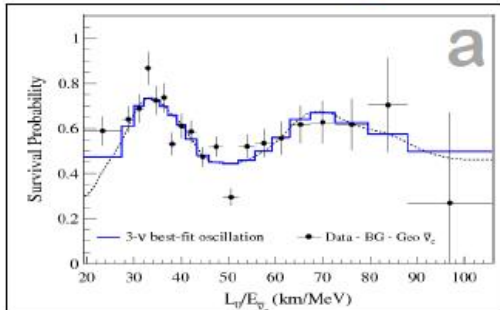


LBL = Long baseline (few x 100 km); **SBL = short baseline** (~1 km)

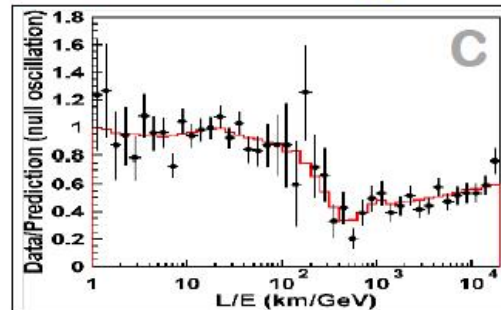
(a) KamLAND reactor [plot]; (b) Borexino [plot], Homestake, Super-K, SAGE, GALLEX/GNO, SNO; (c) Super-K atmosph. [plot], DeepCore, MACRO, MINOS etc.; (d) T2K [plot], NOvA, MINOS, K2K LBL accel.; (e) Daya Bay [plot], RENO, Double Chooz SBL reactor; (f) T2K [plot], MINOS, NOvA LBL accel.; (g) OPERA [plot] LBL accel., Super-K and IC-CD atmospheric.

Now and next... [From E. Lisi (INFN-Bari)]

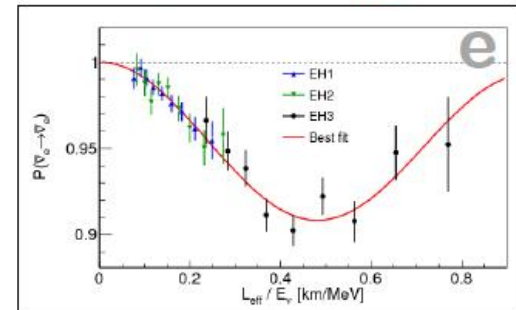
$e \rightarrow e$ ($\delta m^2, \theta_{12}$)



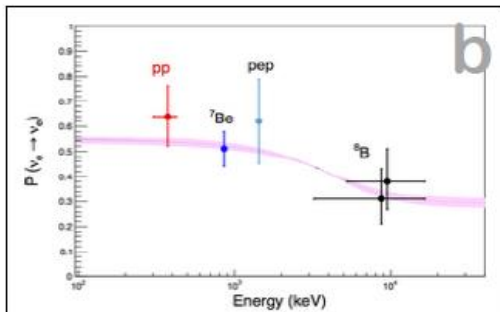
$\mu \rightarrow \mu$ ($\Delta m^2, \theta_{23}$)



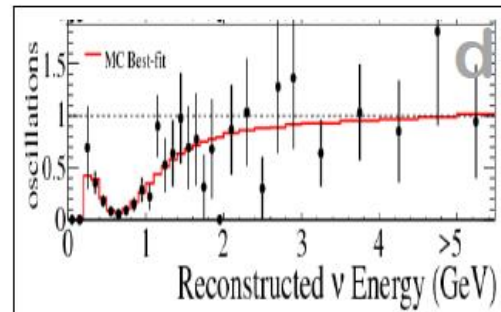
$e \rightarrow e$ ($\Delta m^2, \theta_{13}$)



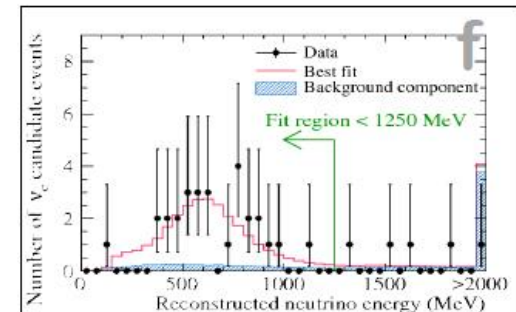
$e \rightarrow e$ ($\delta m^2, \theta_{12}$)



$\mu \rightarrow \mu$ ($\Delta m^2, \theta_{23}$)



$\mu \rightarrow e$ ($\Delta m^2, \theta_{13}, \theta_{23}$)

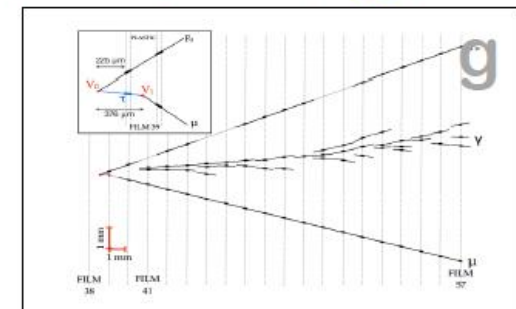


Terra cognita:

δm^2 $|\Delta m^2|$

θ_{12} θ_{23} θ_{13}

$\mu \rightarrow \tau$ ($\Delta m^2, \theta_{23}$)



Terra incognita

5 Knowns:

$$\delta m^2 \sim 7 \times 10^{-5} \text{ eV}^2$$

$$\Delta m^2 \sim 2 \times 10^{-3} \text{ eV}^2$$

$$\sin^2 \theta_{12} \sim 0.3$$

$$\sin^2 \theta_{23} \sim 0.5$$

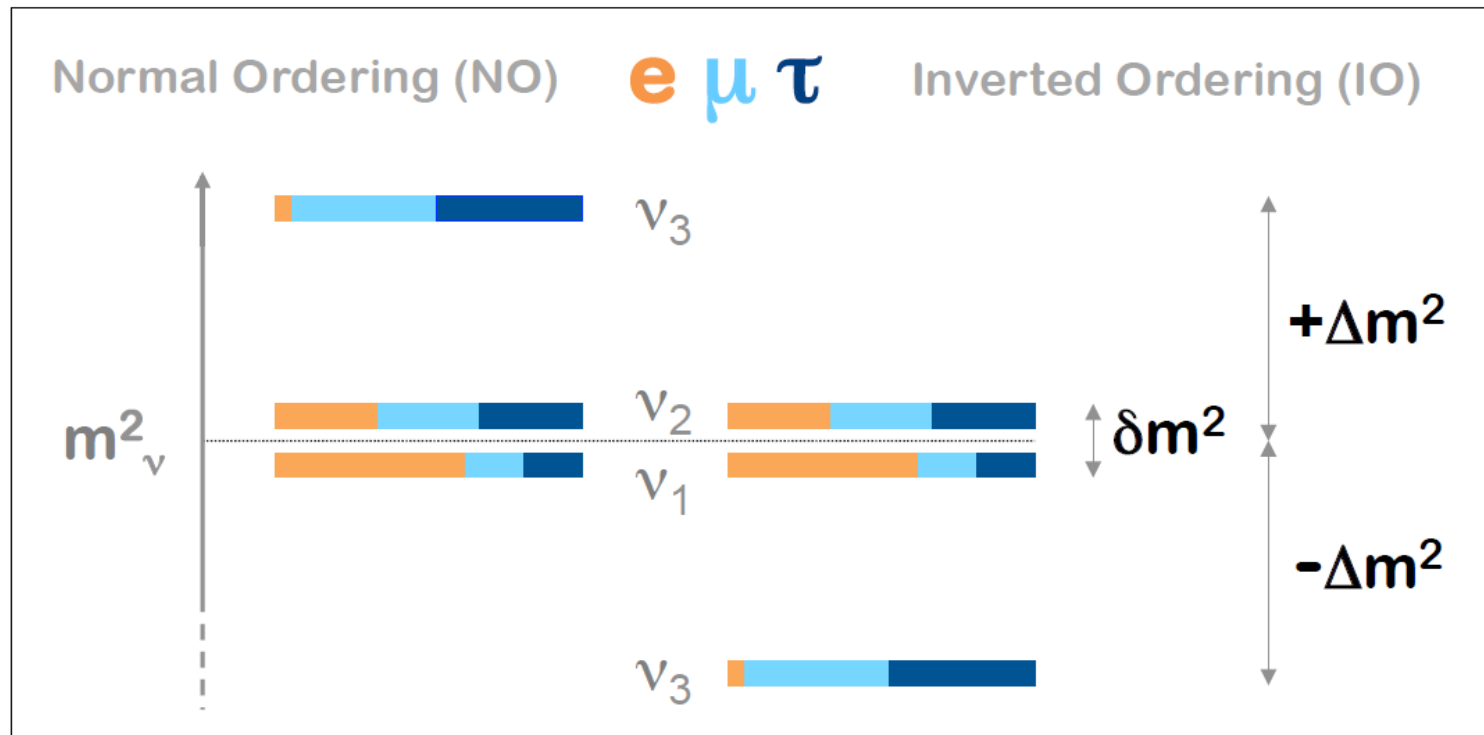
$$\sin^2 \theta_{13} \sim 0.02$$



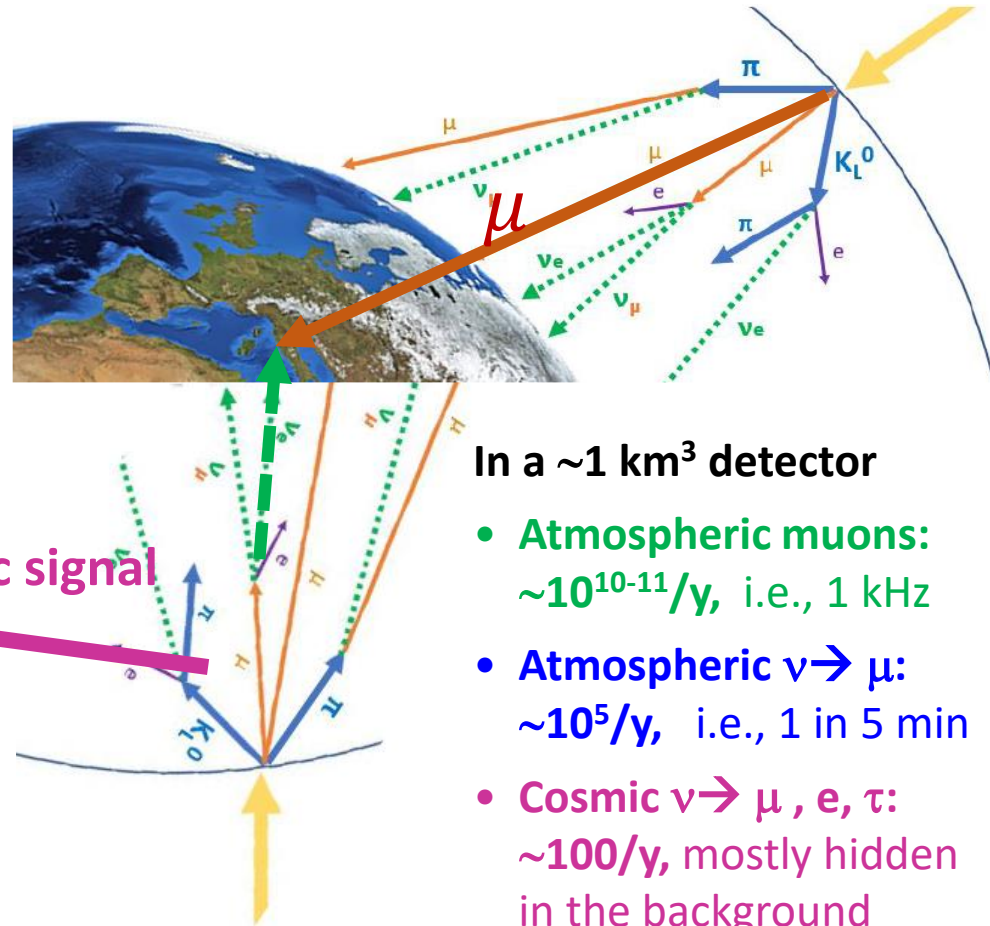
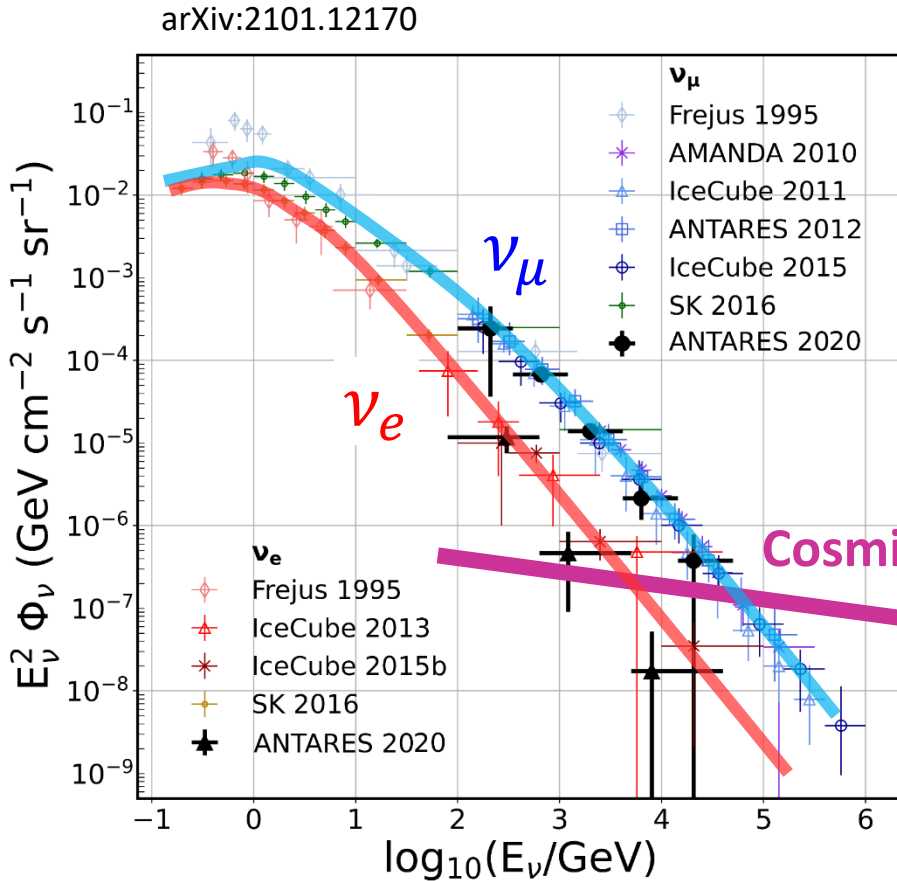
5 Unknowns:

δ = Dirac CPV phase
 $\text{sign}(\Delta m^2)$ = ordering
 octant(θ_{23})

absolute mass scale
 Dirac/Majorana nature



Atmospheric neutrinos: background for cosmic ν



Solution of question 2)

- The event rate of atmospheric neutrinos can be obtained from [Fig. 11.9](#).
- With a threshold of about 1 GeV, the $\Phi_{\nu\mu}(>1 \text{ GeV})=0.65 \text{ cm}^{-2} \text{ s}^{-1}$, and that of $\Phi_{\nu e}(>1 \text{ GeV})$ a factor of two lower. The neutrino must interact inside the detector to produce a visible event.
- The event rate for $T = 1 \text{ y}$ and for the flavor $i = \mu; e$ is given by

$$N_{\nu_i} = \Phi_{\nu_i}(>1 \text{ GeV}) \cdot \sigma_{\nu} \cdot N_T \cdot T ,$$

- $\sigma_{\nu} = \sigma_0 E$ is the neutrino cross-section. After averaging over neutrinos and antineutrinos, we have $\sigma_0 \approx 0.5 \cdot 10^{-38} \text{ cm}^2/\text{GeV}$
- $N_T = 6 \times 10^{32} \text{ nucleons/ton}$ is the number of target nucleons in 1 ton,
- $T = 3.15 \cdot 10^7 \text{ s}$ is the number of seconds in 1 year.
- Thus, inserting the numerical values, we obtain

$$N_{\nu_{\mu}} = 66 \text{ kton/y} \quad ; \quad N_{\nu_e} = 34 \text{ kton/y} .$$

Solution of question 3)

- At low energy, neutrinos exceed antineutrinos due to the fact that CR protons are more abundant than neutrons
- Above few hundreds GeV, neutrinos from K decay are more abundant than from pions. Thus, **conservation of the strangeness (S) and baryon (B) quantum numbers** are responsible for the difference
- Consider the production of charged kaons from pp interactions:
 - K^+ (B=0, S=1) is produced in association with Λ (B=1, S=-1);
 - K^- (B=0, S=-1) requires at least one associated baryon (B=1) **and** an additional strange meson (S=1).
- K^+ are generated much more frequently than K^- because of the associated production with the Λ .

Solution of question 5)

- In water ($n=1.33$) for $z=1$ and $\beta=1$ the number of photons for energy interval and unit path length is:

$$\frac{d^2 N}{dx dE} = \frac{z^2 \alpha}{\hbar c} \sin^2 \mathcal{G}_c.$$

- With: $\cos \mathcal{G}_c = (1/n) = 0.75 \rightarrow \mathcal{G}_c = 42^\circ$
- The $\Delta\lambda=(600-300)$ nm corresponds to $\Delta E=2$ eV, thus

$$\frac{d^2 N}{dx} = \frac{(1/137)}{(197 \text{MeV fm})} \cdot \Delta E \cdot \sin^2 \mathcal{G}_c = \frac{2 \text{eV} \times (0.67)^2}{27 \cdot 10^3 \times 10^6 \text{eV} \times 10^{-13} \text{cm}} = \frac{330}{\text{cm}}$$